



“Tidal effects in multi-planetary systems”

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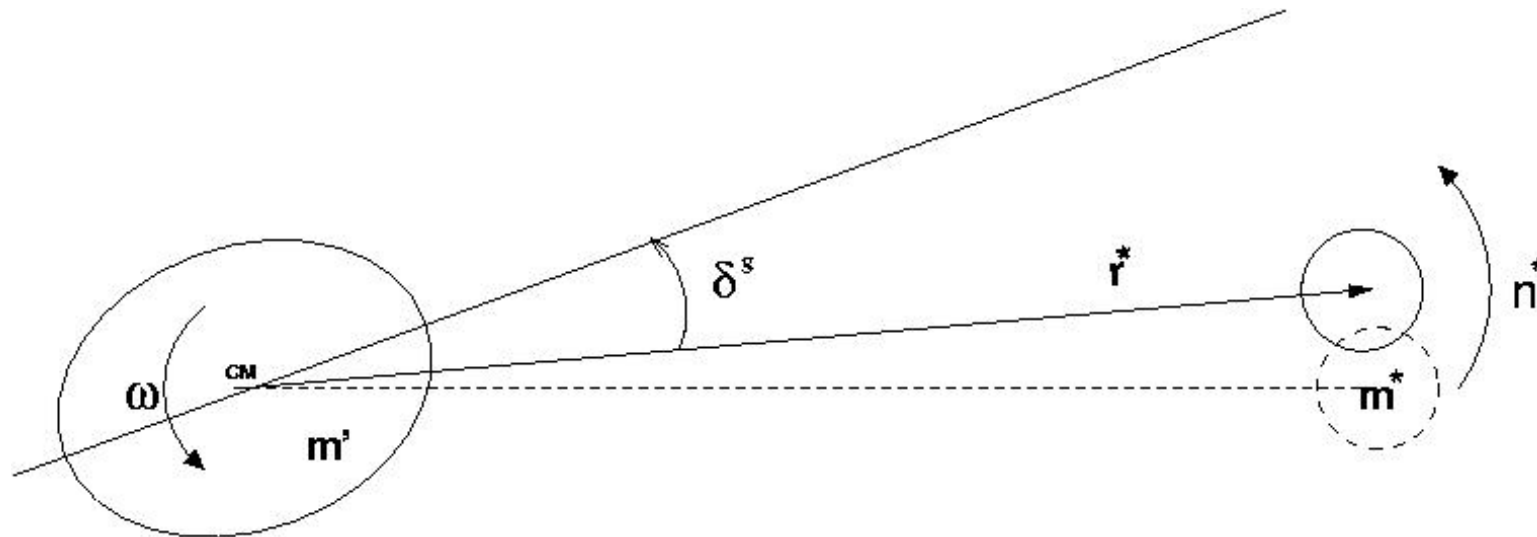
FUNDP, Namur

8 mai 2012

Tidal effect



Tidal effect



$$P_{\text{orb}} = P_{\text{rot}}$$

perpendicular axis ($\varepsilon=0$)

circular orbits ($e=0$)

Case of the Moon and
remaining satellites of
the Solar System

Global Picture:

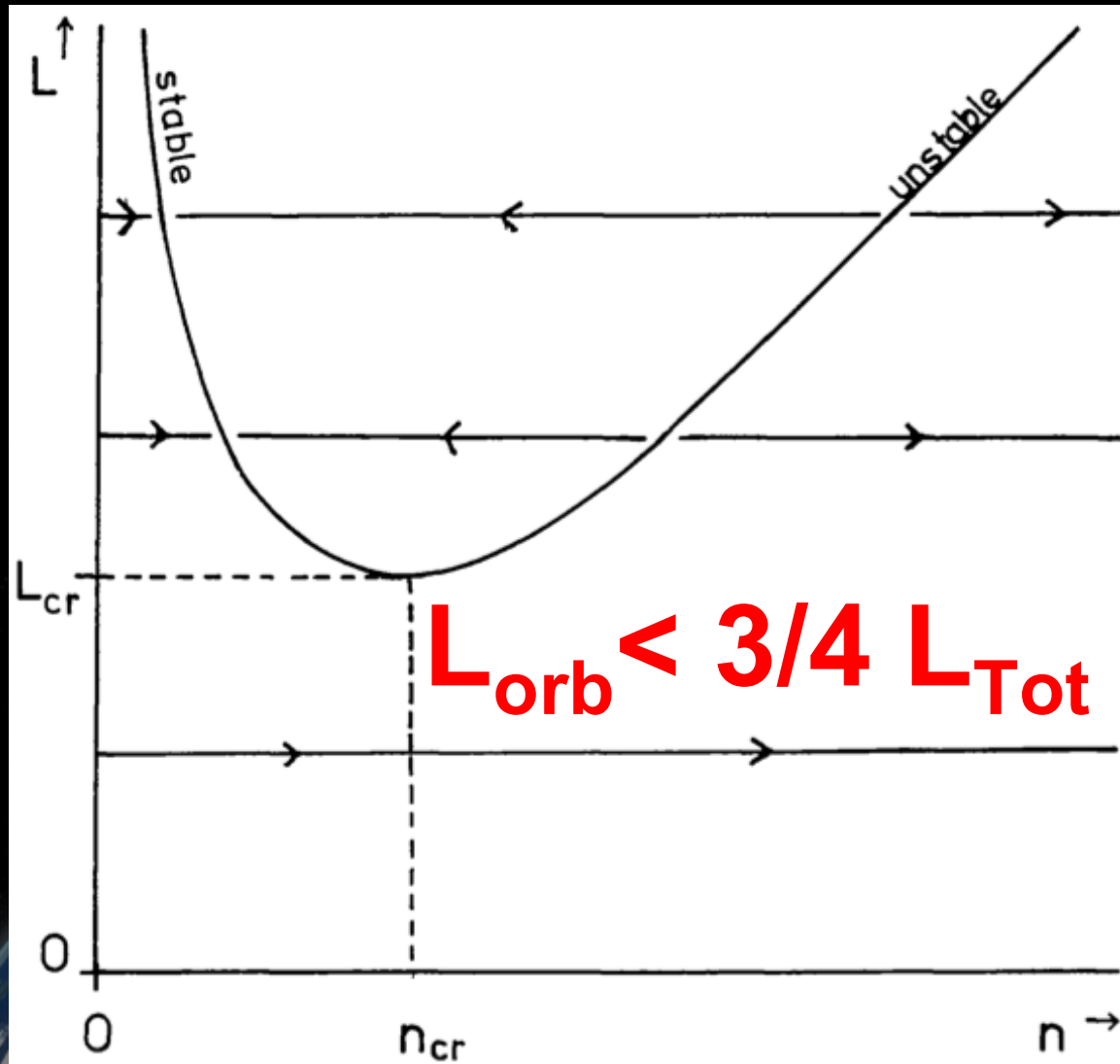
Poincaré (1898)

Landau & Lifshitz (1960)

Hut (1980)

$P_{\text{orb}} = P_{\text{rot}}$
perpendicular
axis ($\varepsilon=0$)

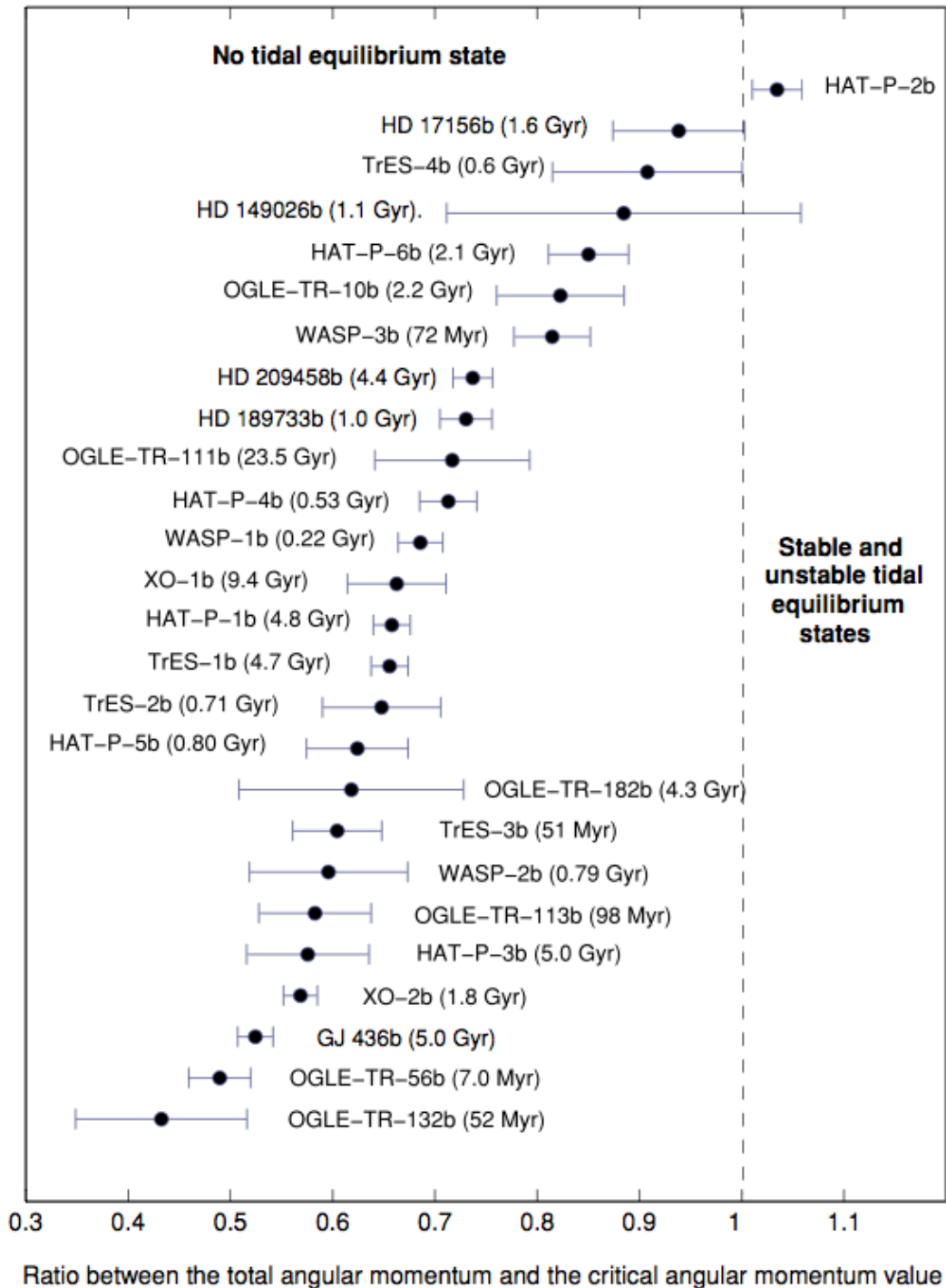
circular
orbits ($e=0$)



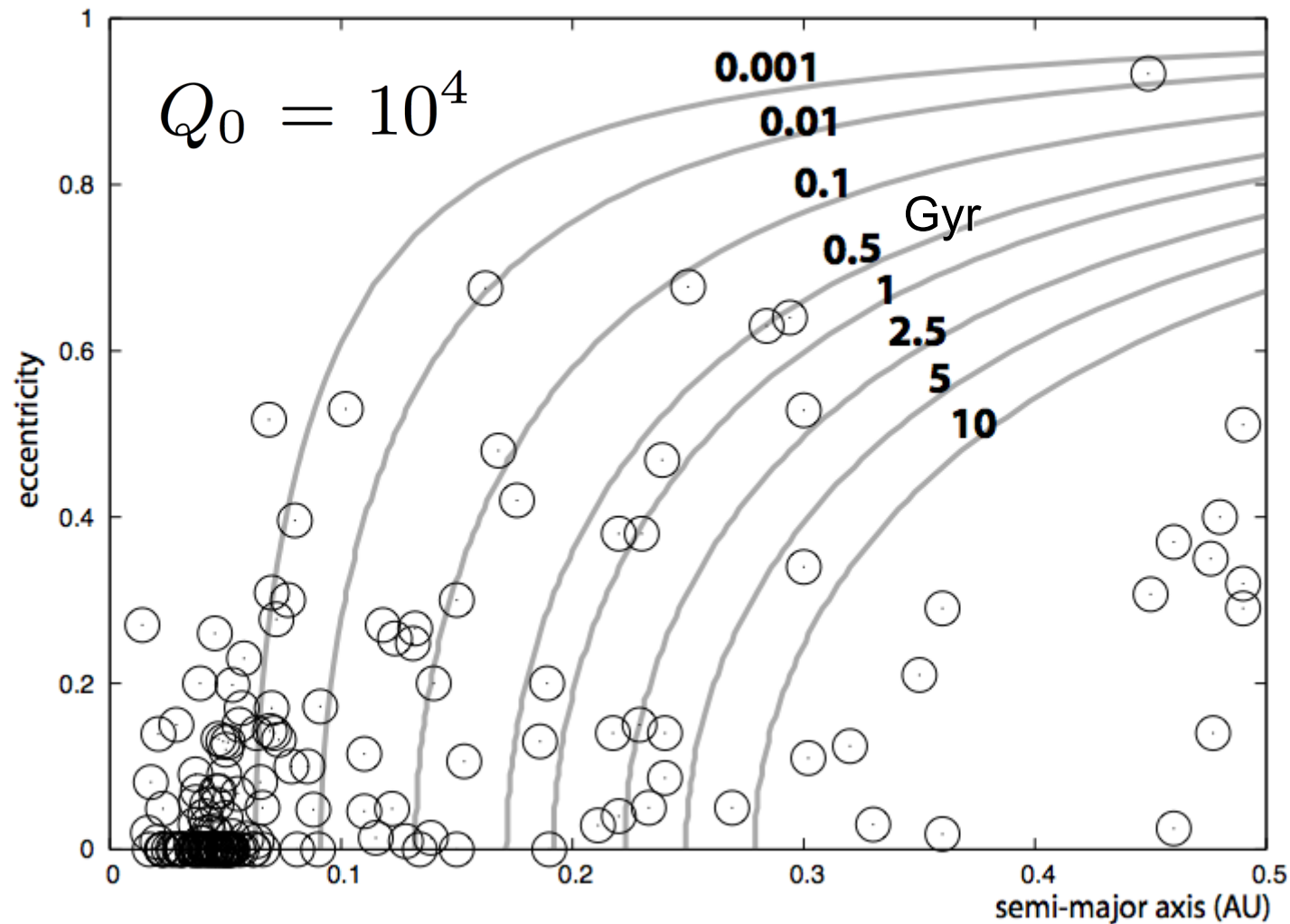
$$L_{\text{cr}} = 4 \left\{ \frac{1}{27} G \frac{M^3 m^3}{M + m} (I_1 + I_2) \right\}^{1/4}$$

Falling Hot-Jupiters

Levrard *et al.* (2009)

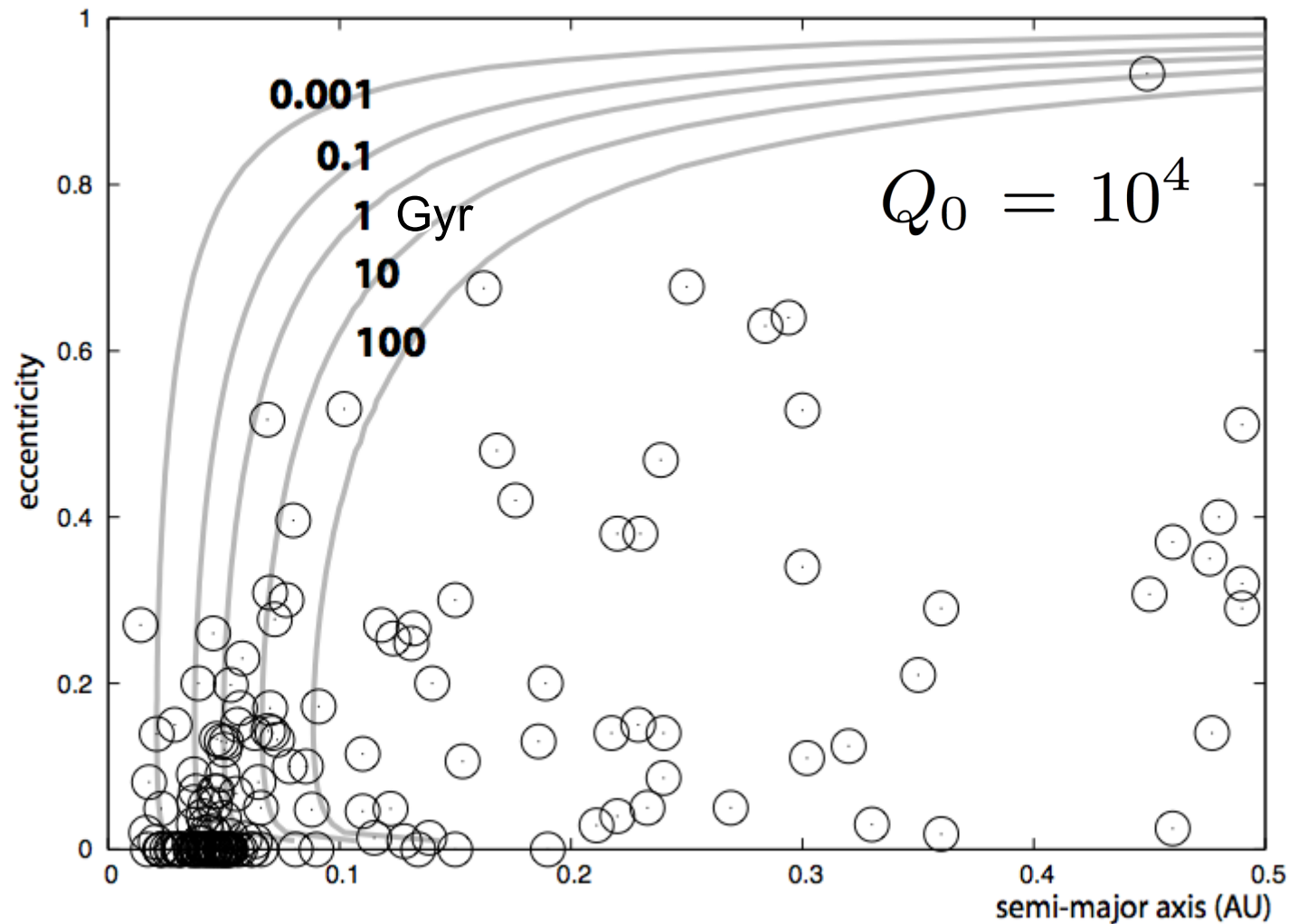


Spin evolution time-scale:



Correia & Laskar, *Exoplanets* (2010)

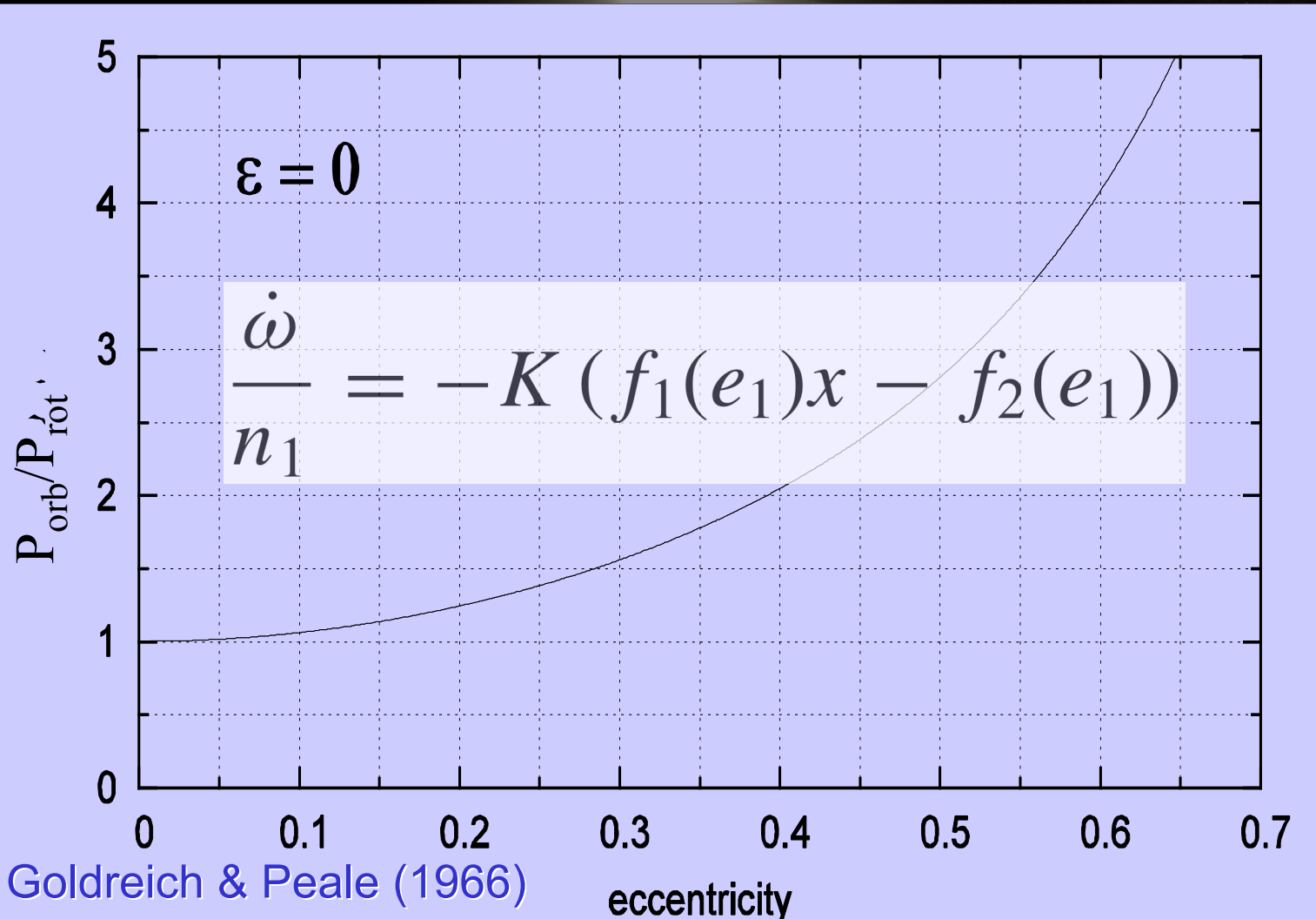
Orbital evolution time-scale:



Correia & Laskar, *Exoplanets* (2010)

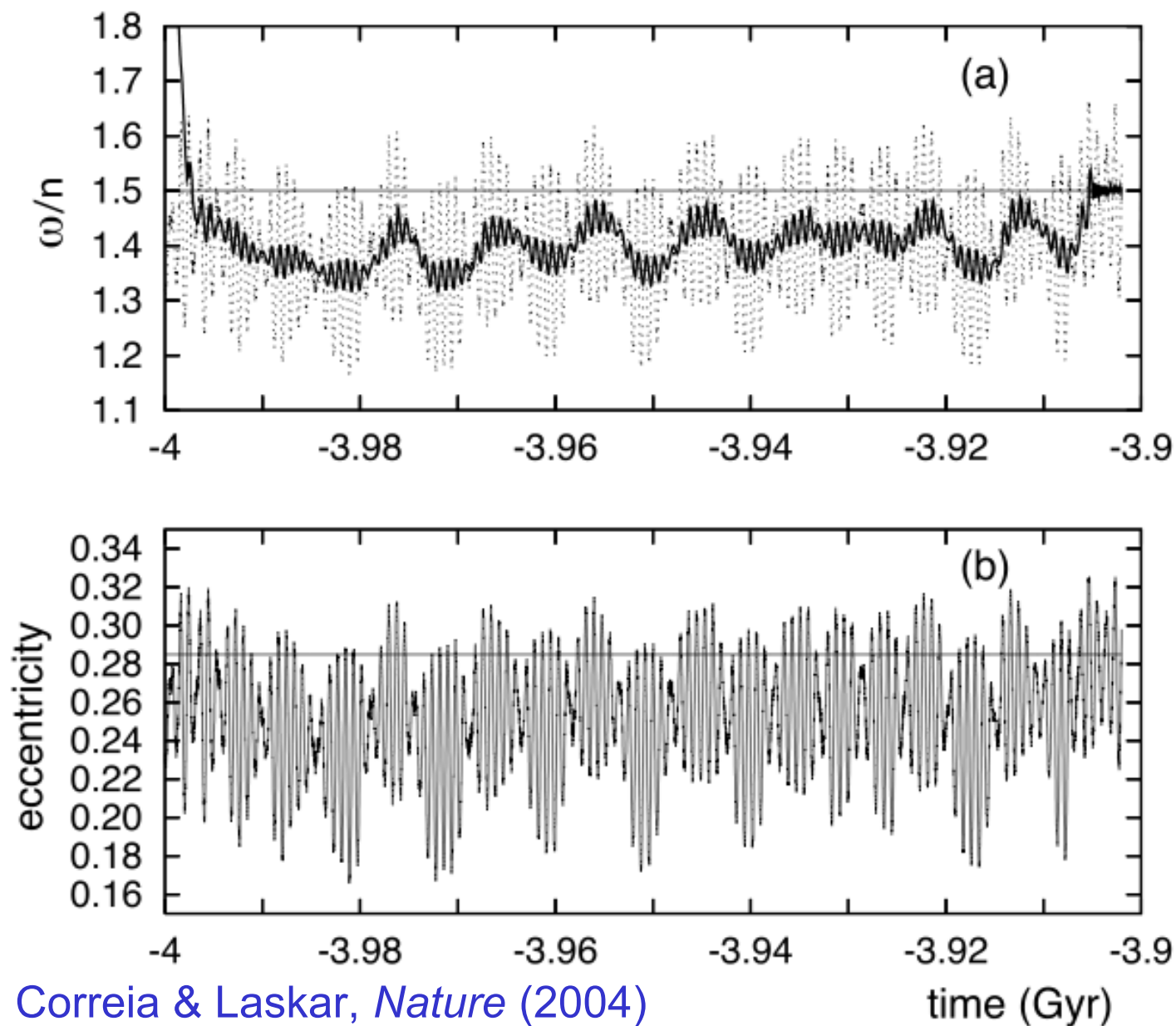
Intermediate evolution:

$$\omega_{eq} / n = f_2/f_1(e)$$



Rotation tidal wobble

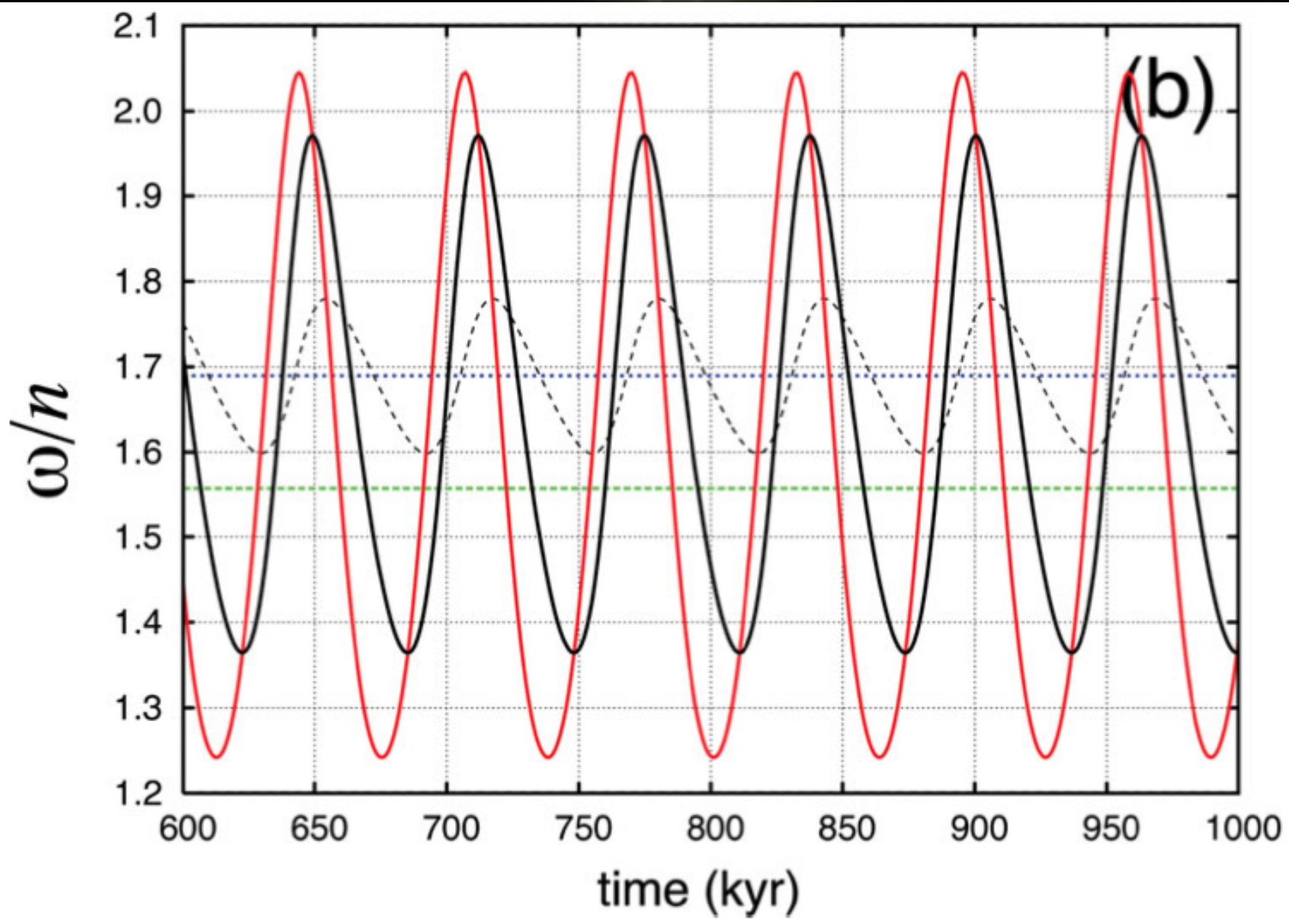
Mercury



Correia & Laskar, *Nature* (2004)

Rotation tidal wobble

Jupiter-size planets



linearized system $x = \omega / n$

$$\delta \dot{e}_1 = -A \sin \varpi,$$

$$\delta e_1 = \Delta e \cos(gt + \varpi_0)$$

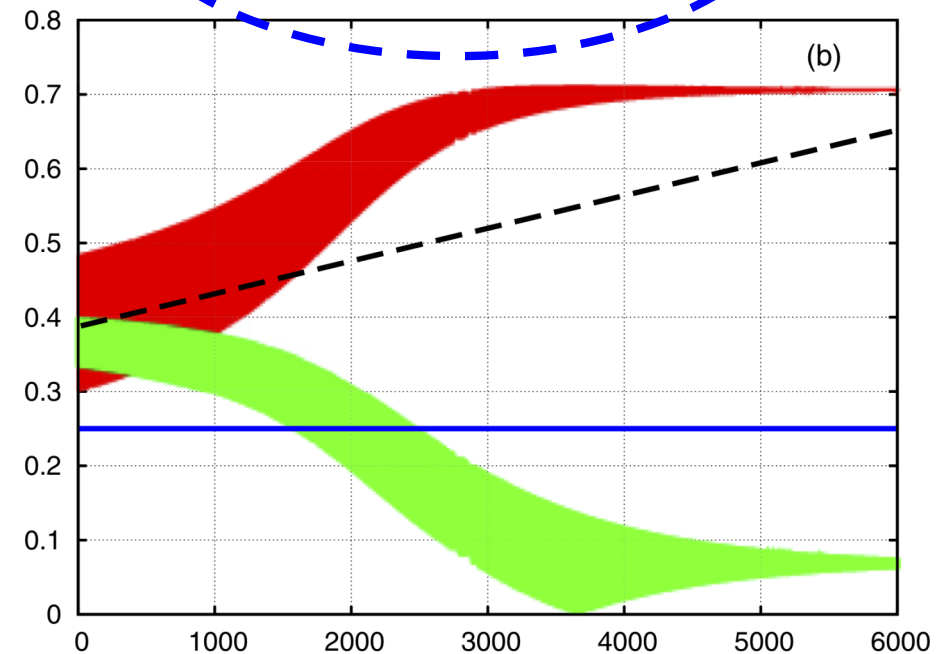
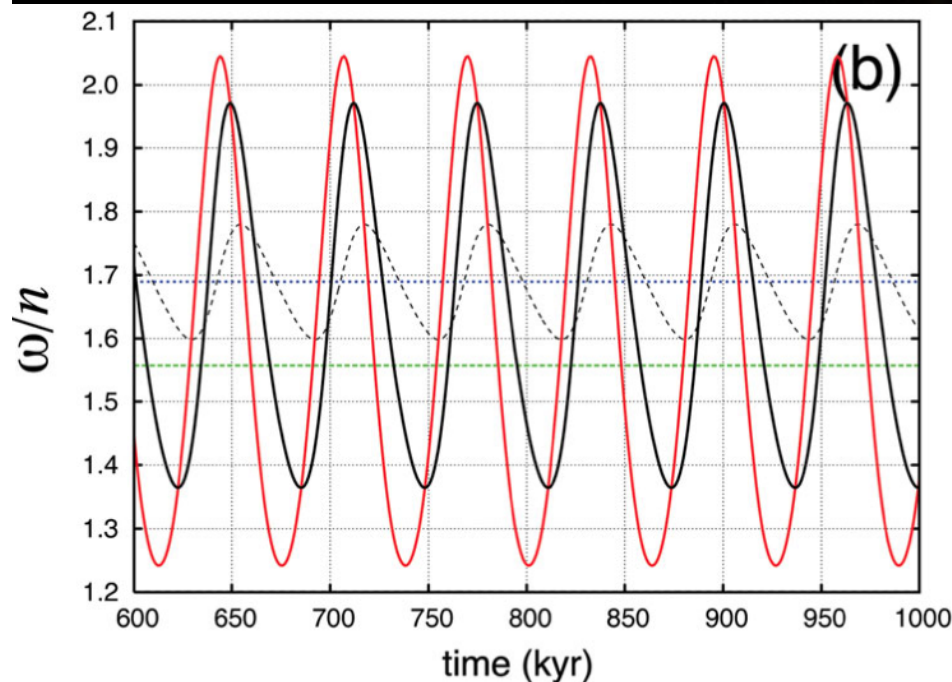
$$\dot{\varpi} = g$$

$$\delta x = \Delta x \cos(gt + \varpi_0 - \phi)$$

$$\delta \dot{x} = -\nu_x \delta x + \nu_e \delta e_1,$$

drift on the eccentricity

$$\dot{e}_1 = -A \sin(gt + \varpi_0) - \frac{g_x A}{2g} \Delta x \sin(2gt + 2\varpi_0 - \phi) - \frac{g_e A}{2g} \Delta e \sin(2gt + 2\varpi_0) + \frac{g_x A}{2g} \Delta x \sin \phi.$$



moderate close-in planets

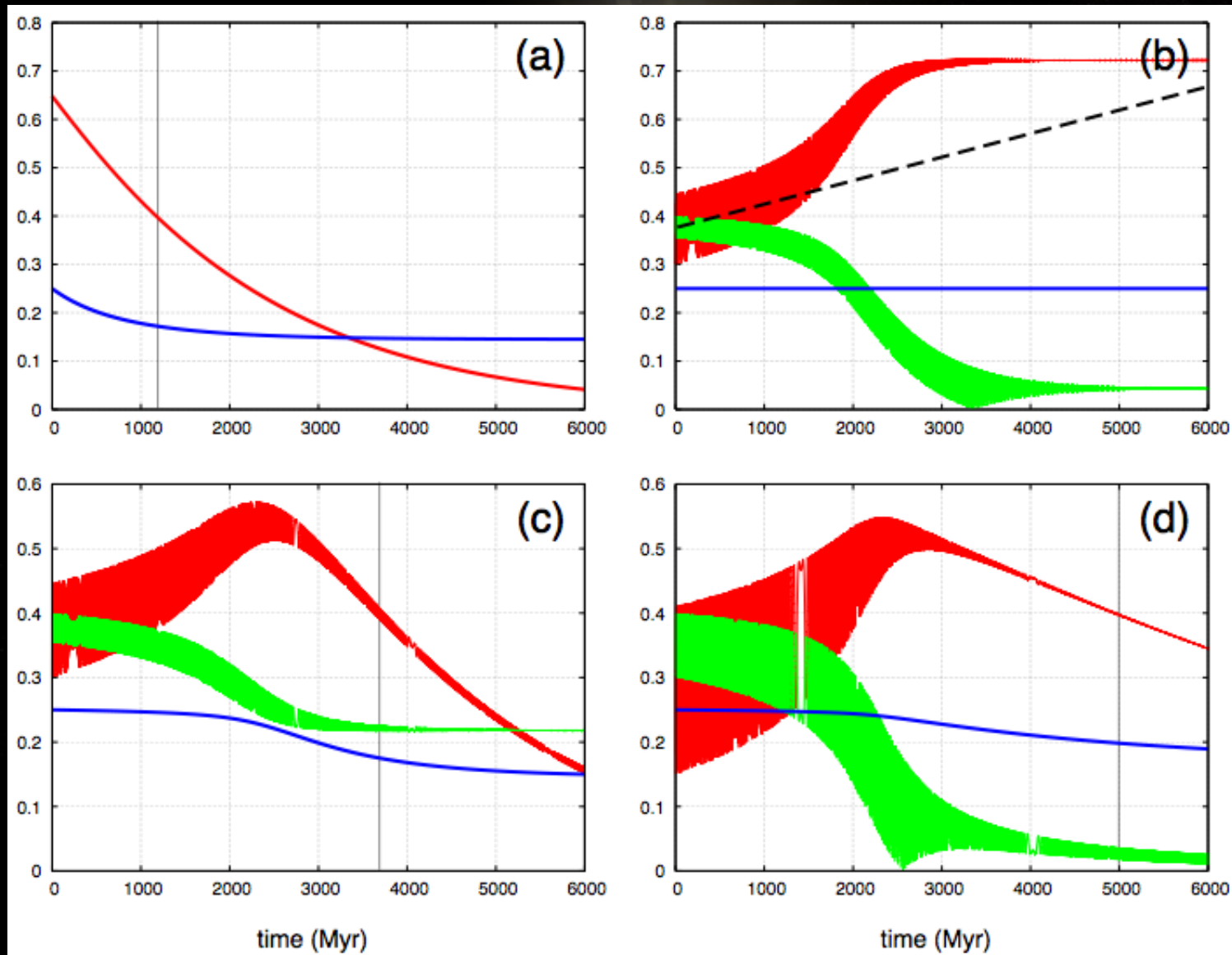
Table 1

Single Planetary Systems with $0.1 < a_1 < 0.3$ and $e_1 > 0.3$

Star (Name)	a_1 (AU)	e_1	m_1 (M_J)	m_0 ($M_\odot$)	Age (Gyr)	τ (Gyr)
HD 108147	0.102	0.53	0.26	1.19	2.0	0.01
CoRoT-10	0.105	0.53	2.75	0.89	3.0	0.24
HD 33283	0.145	0.48	0.33	1.24	3.2	0.34
HD 17156	0.163	0.68	3.19	1.28	3.4	0.44
HIP 57050	0.164	0.31	0.30	0.34	...	39.4
HD 117618	0.176	0.42	0.18	1.05	3.9	2.06
HD 45652	0.228	0.38	0.47	0.83	...	93.3
HD 90156	0.250	0.31	0.06	0.84	4.4	35.8
HD 37605	0.260	0.74	2.84	0.80	10.7	10.6
HD 3651	0.284	0.63	0.20	0.79	5.1	15.5

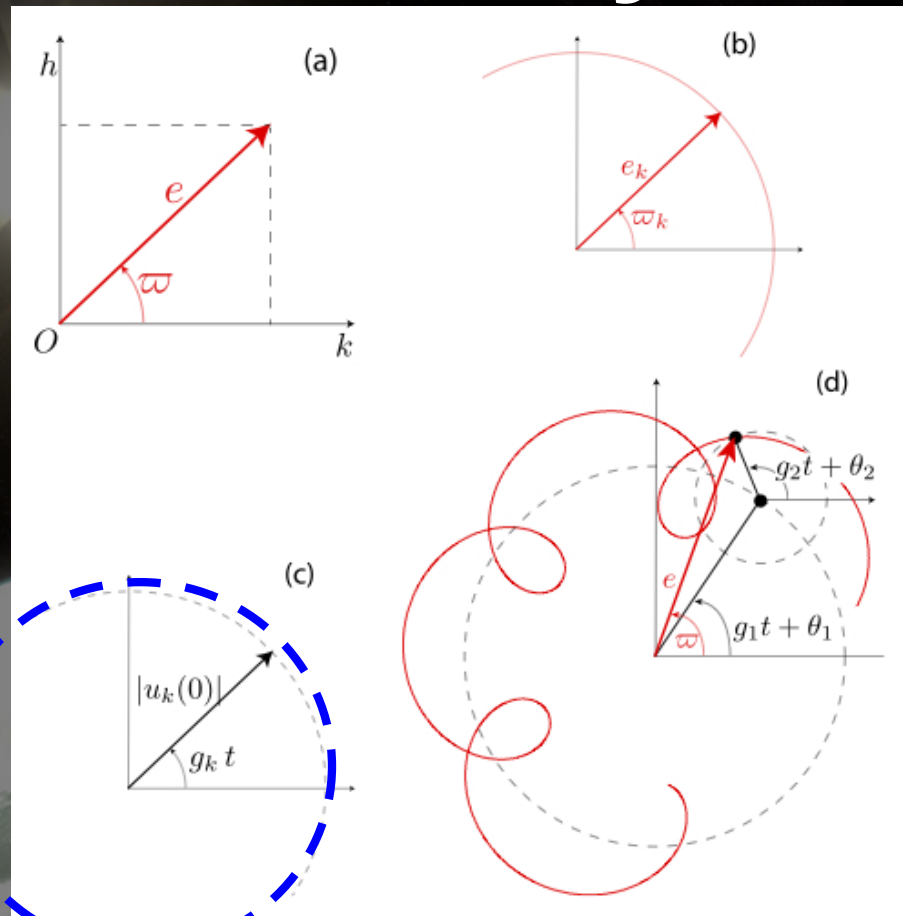
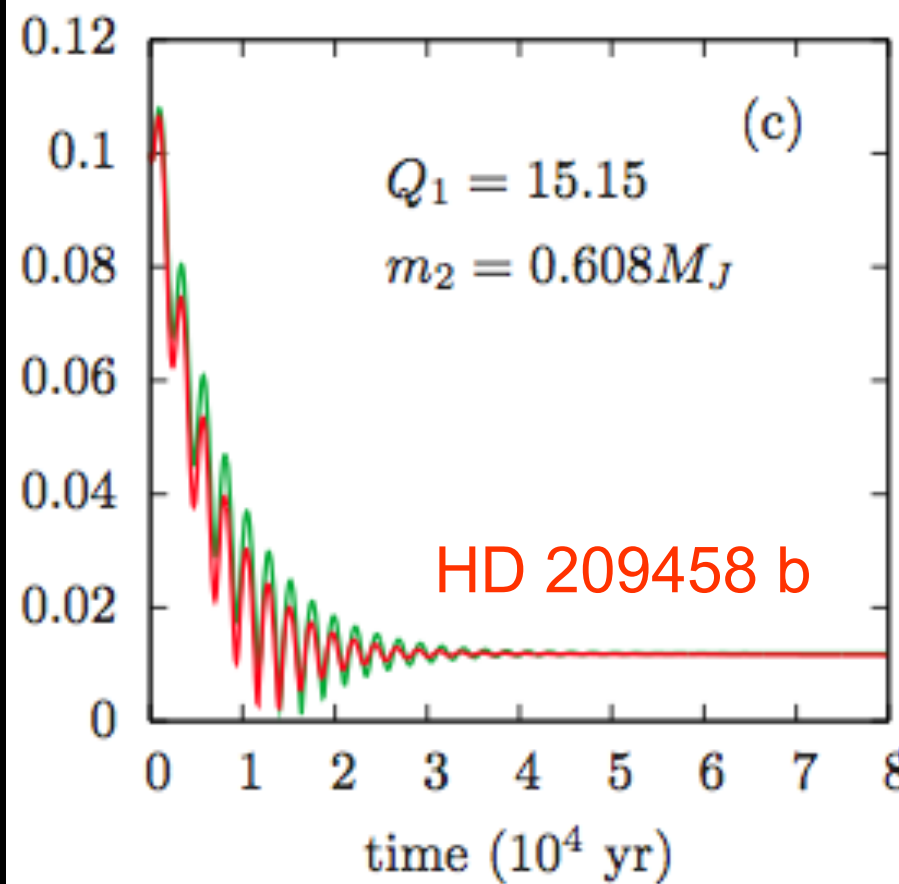
eccentricity pumping

HD 117618 b



Correia, Boué & Laskar, *ApJ Lett* (2012)

equilibrium eccentricity



$$z_k = e_k e^{i\varpi_k}, k = 1, \dots, n$$

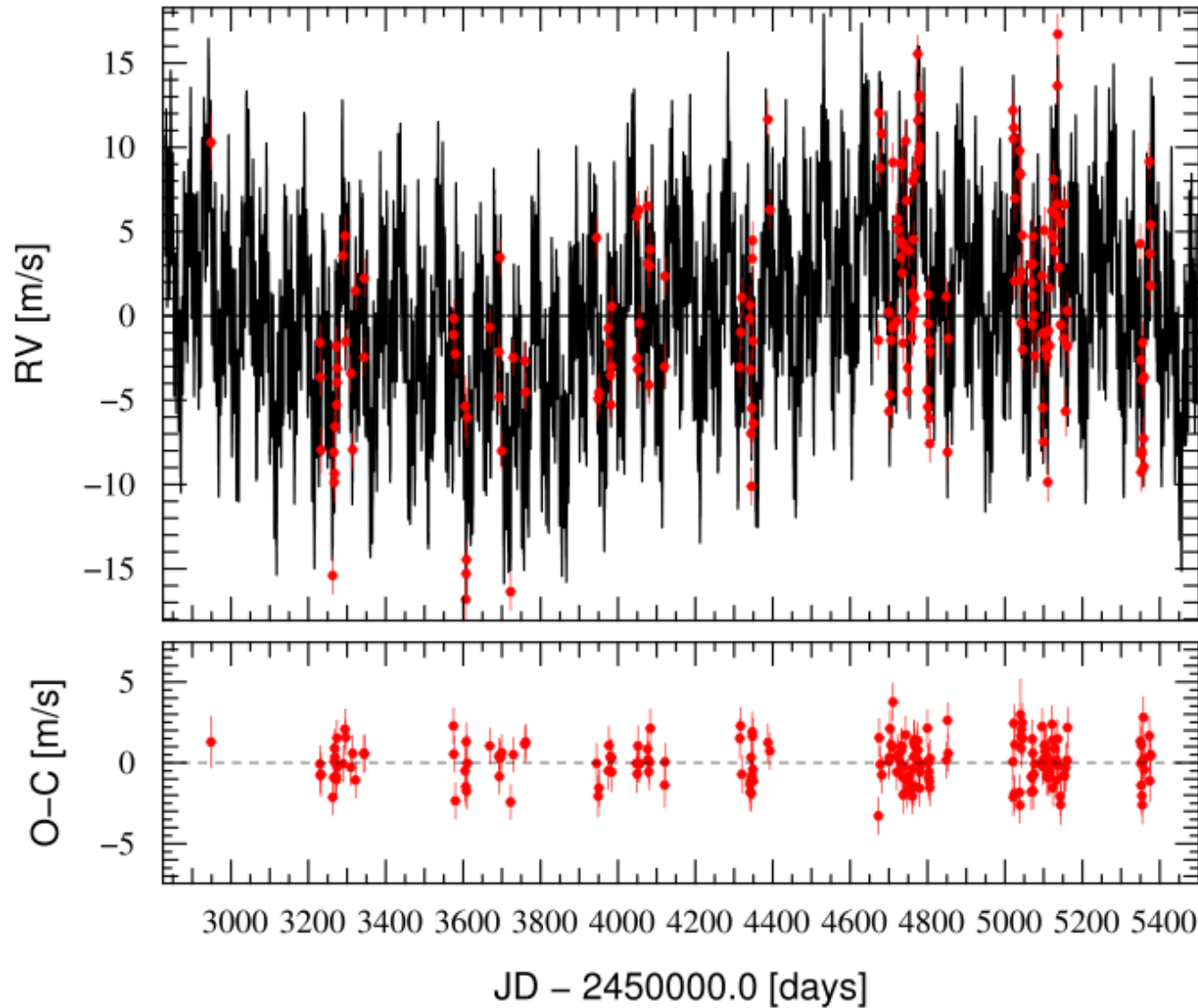
$$u_k(t) = u_k(0) e^{-\gamma_k t} e^{ig_k t}$$

$$z_1(t) = \cancel{S_{11}u_1(t)} + S_{12}u_2(t)$$

$$z_2(t) = \cancel{S_{21}u_1(t)} + S_{22}u_2(t)$$

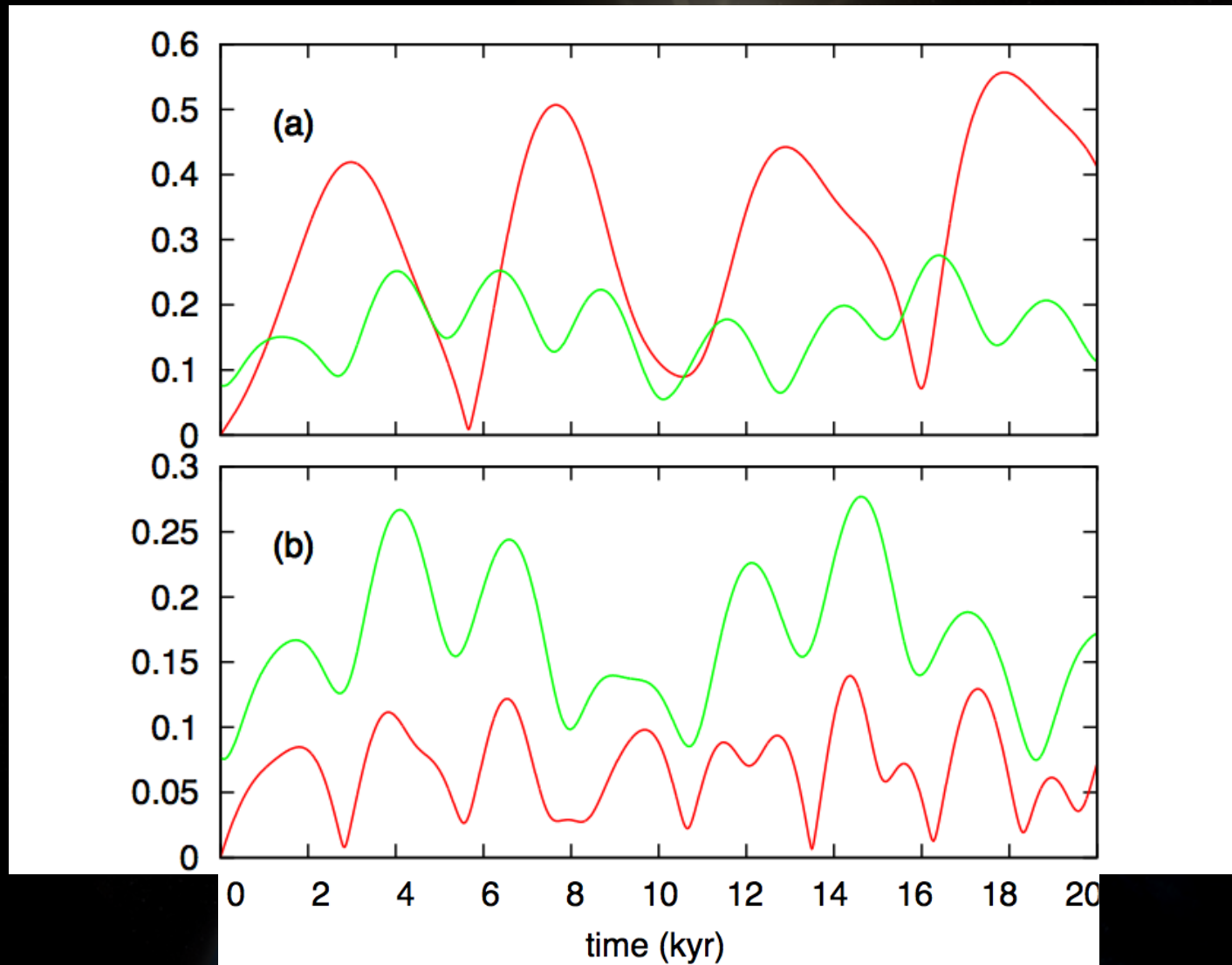
Mardling, *MNRAS* (2007); Laskar, Boué & Correia, *A&A* (2012)

HD 10180, seven planets! (or nine)



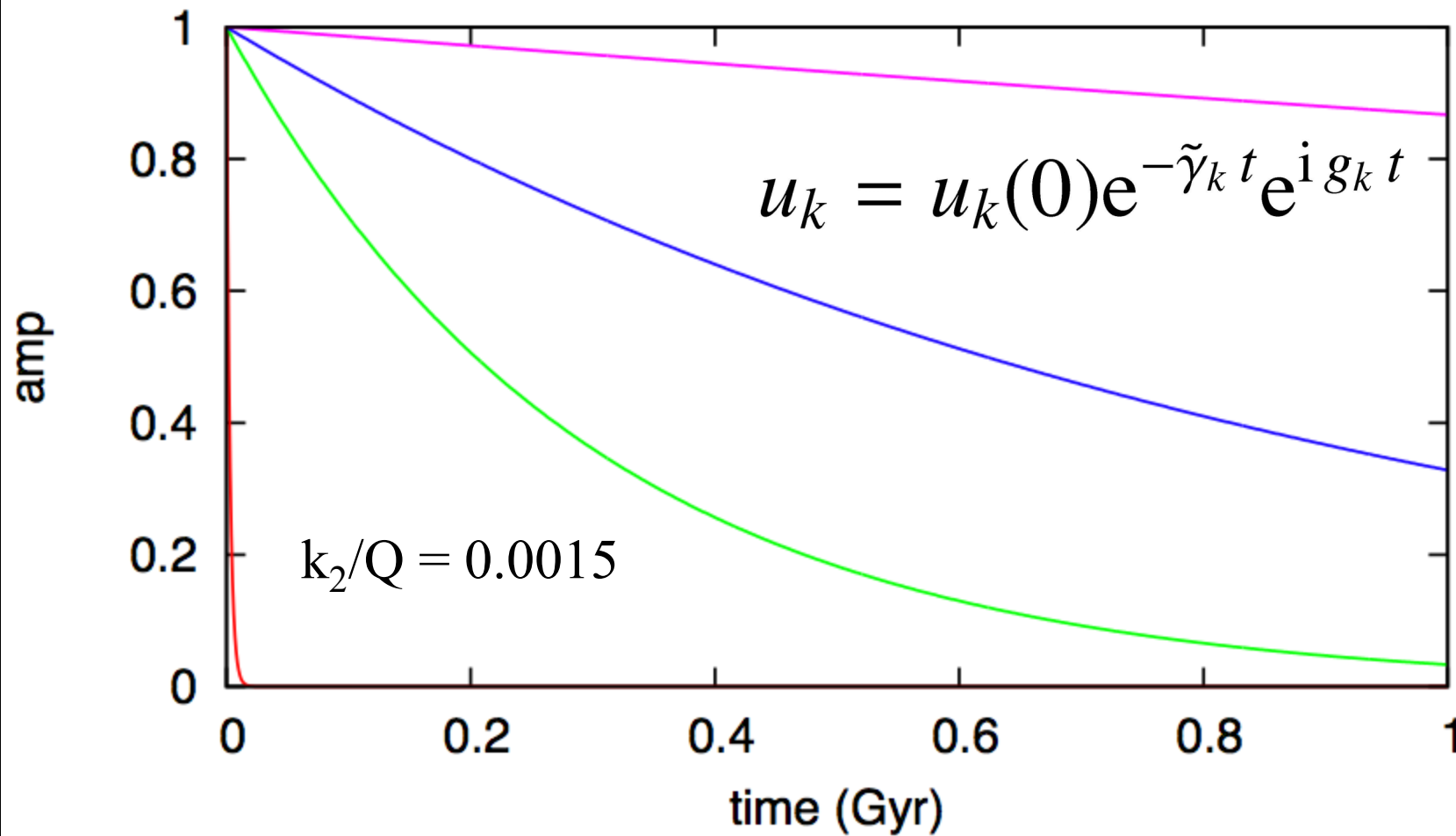
Lovis et al. (A&A 2011)

unstable system?



Lovis et al. (A&A 2011)

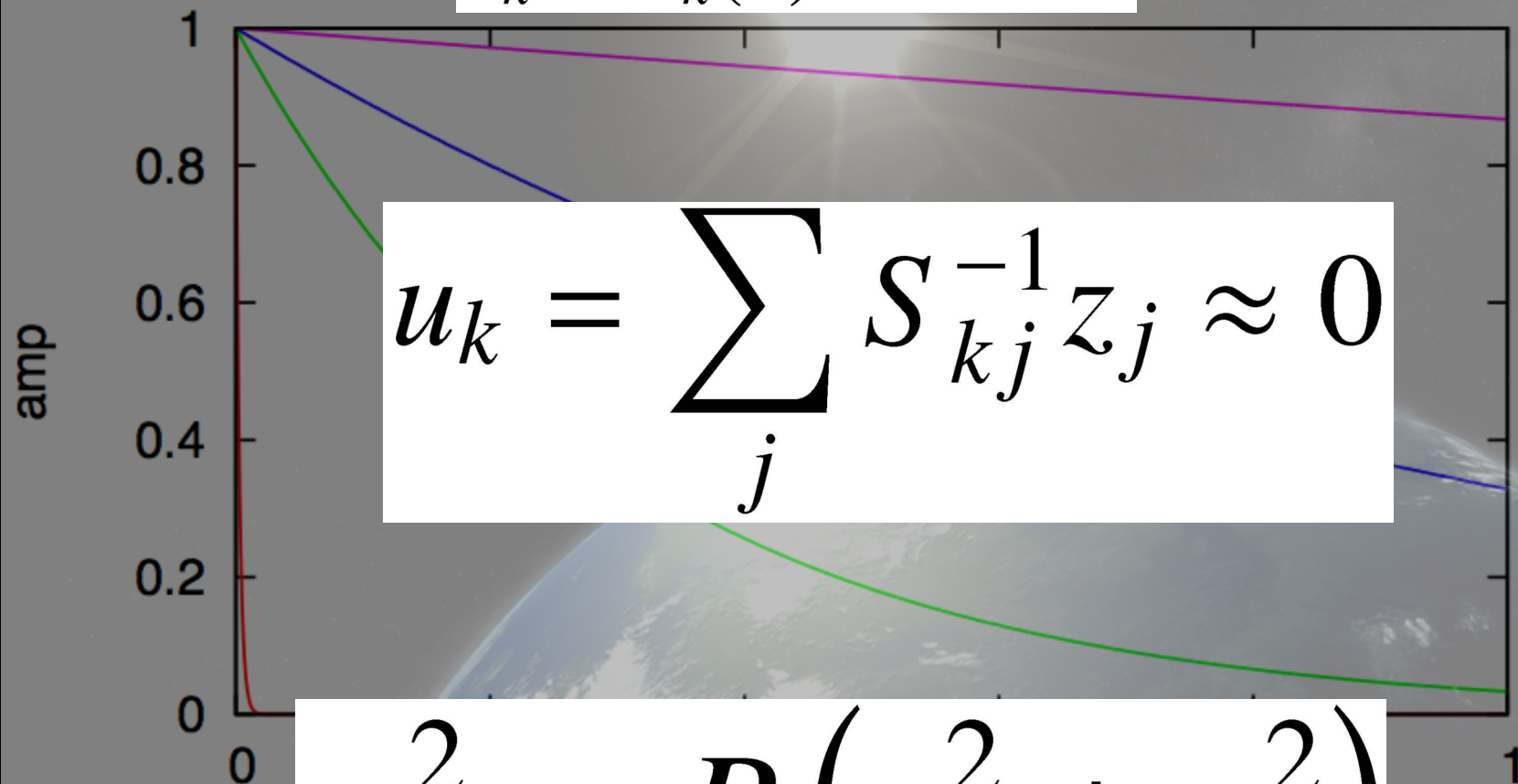
tidal dissipation



Lovis et al. (A&A 2011)

tidal constraint

$$u_k = u_k(0)e^{-\tilde{\gamma}_k t} e^{i g_k t}$$

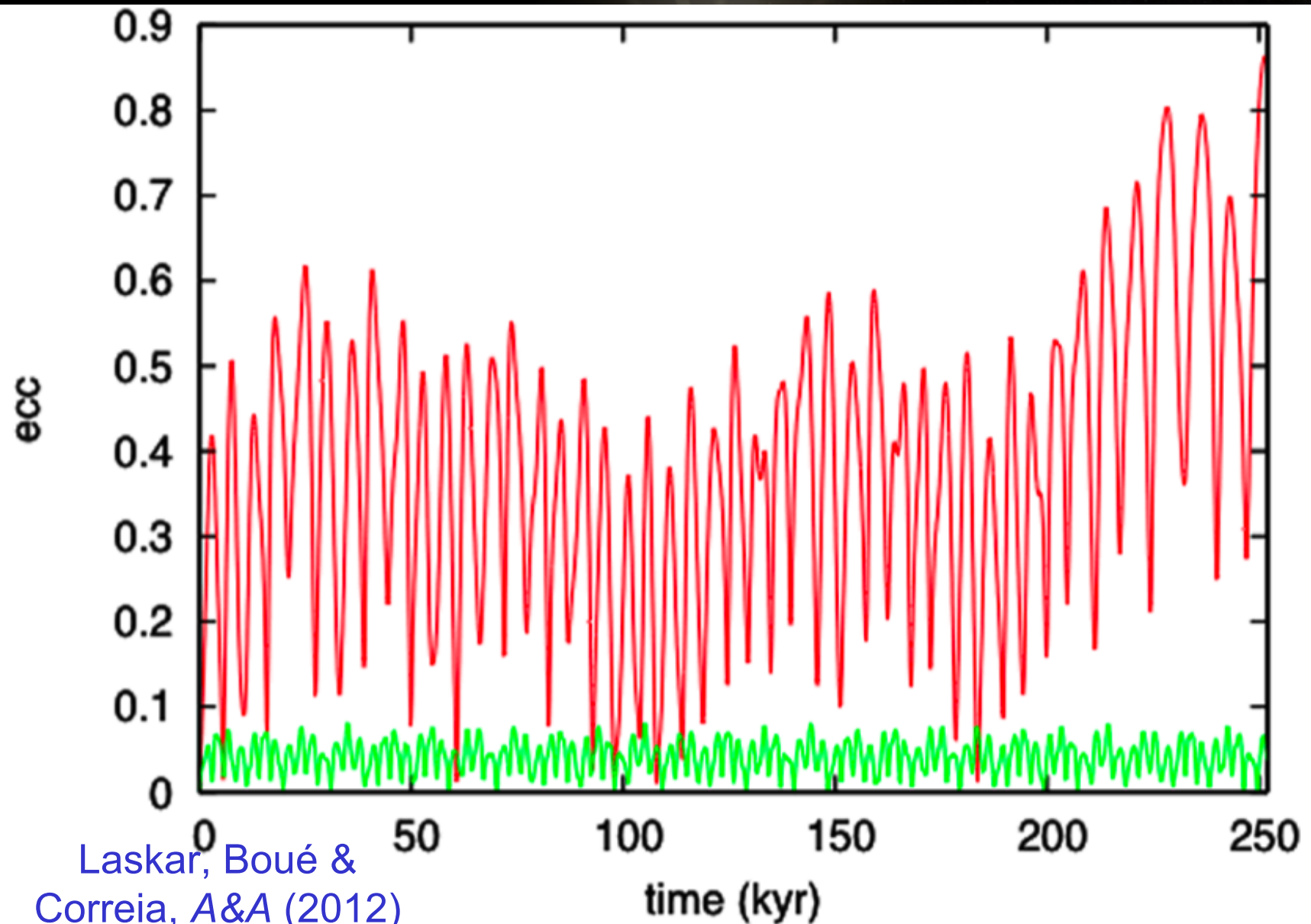


$$u_k = \sum_j S_{kj}^{-1} z_j \approx 0$$

$$\chi_R^2 = R \left(u_1^2 + u_2^2 \right)$$

Lovis et al. (A&A 2011)

eccentricity evolution

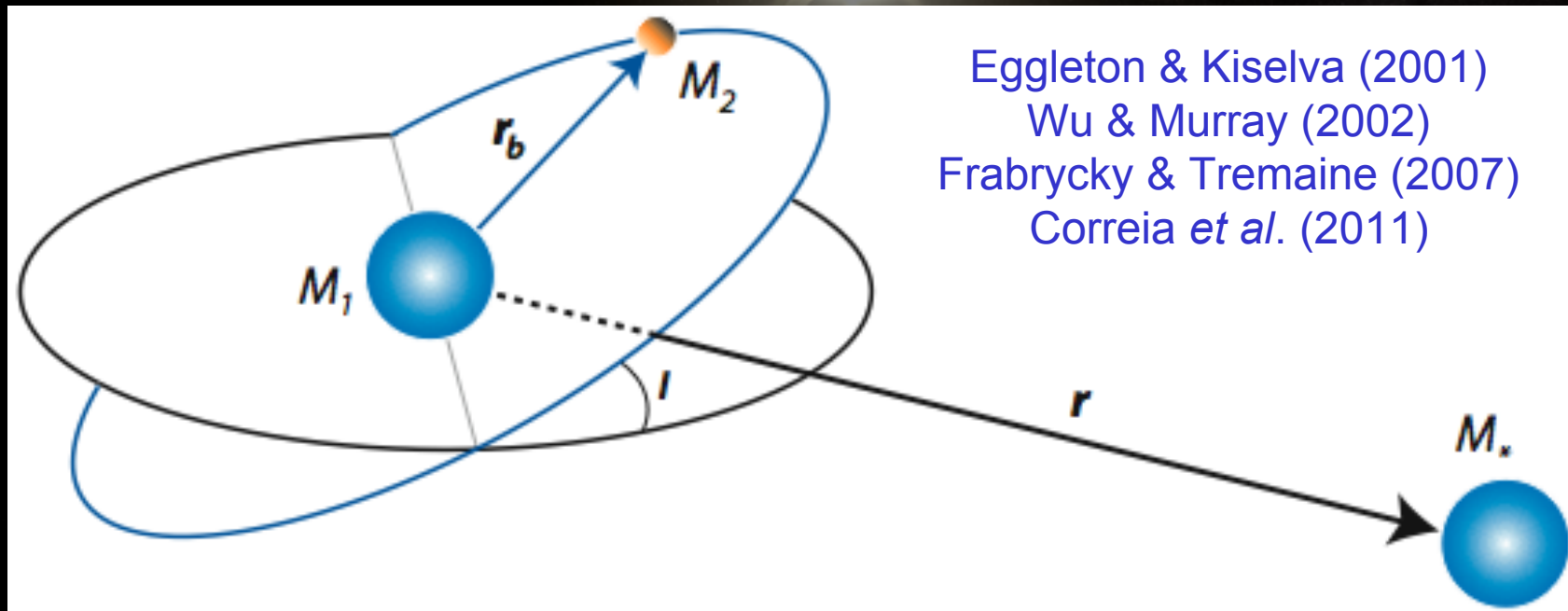


HD 80606

(inner restricted problem)

Naef et al, 2001

($a_p = 0.45$ AU, $e_p = 0.92$)

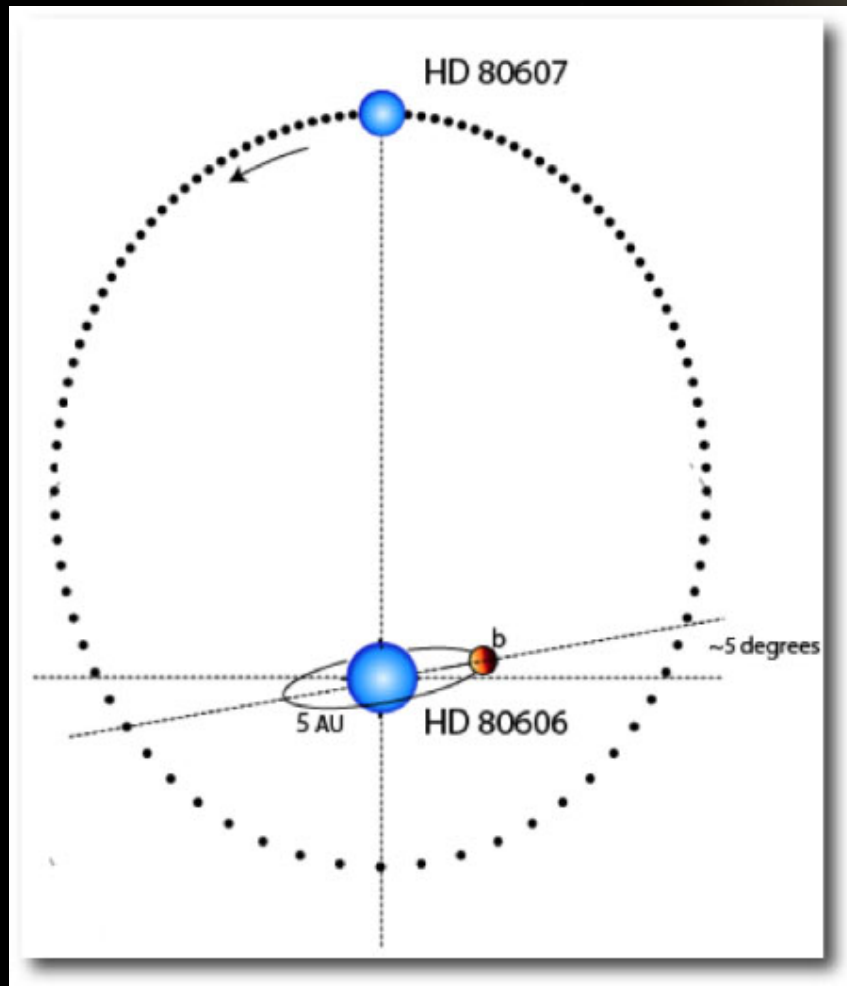


Kozai effect:
($I > 39^\circ$)

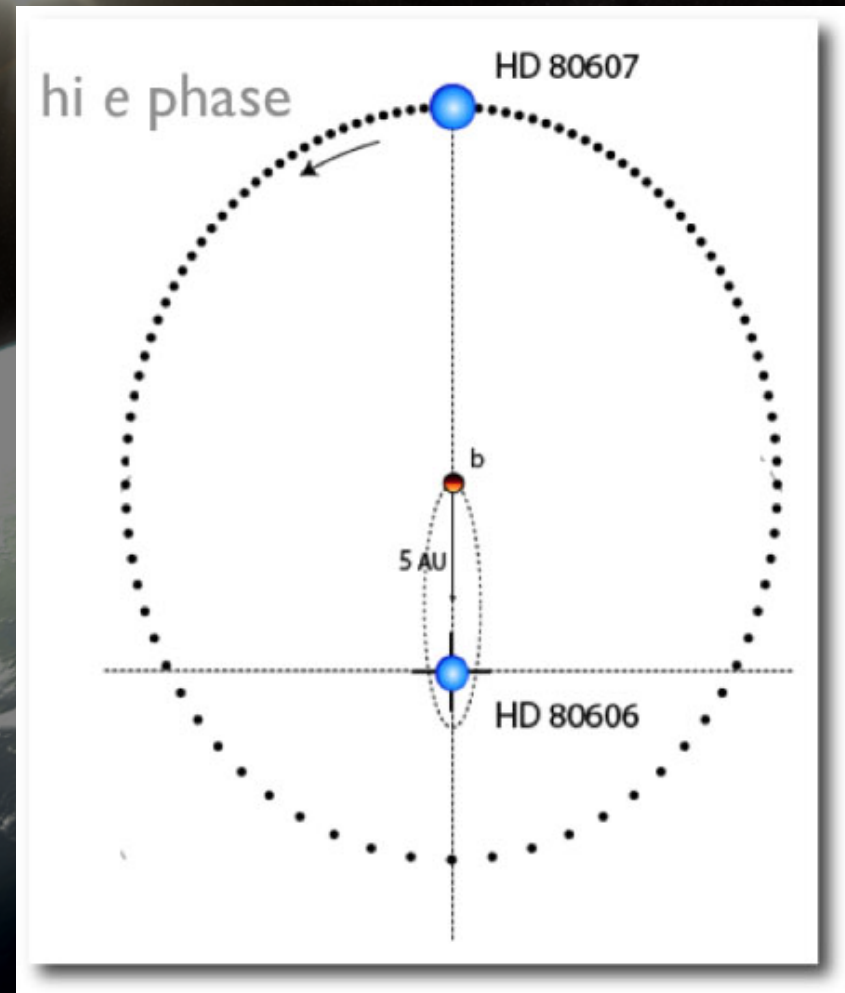
$$\sqrt{1 - e_1^2} \cos I = h_1 = Cte$$

HD 80606

$i \sim 85^\circ$, $e \sim 0$



$i \sim 45^\circ$, $e \sim 0.99$

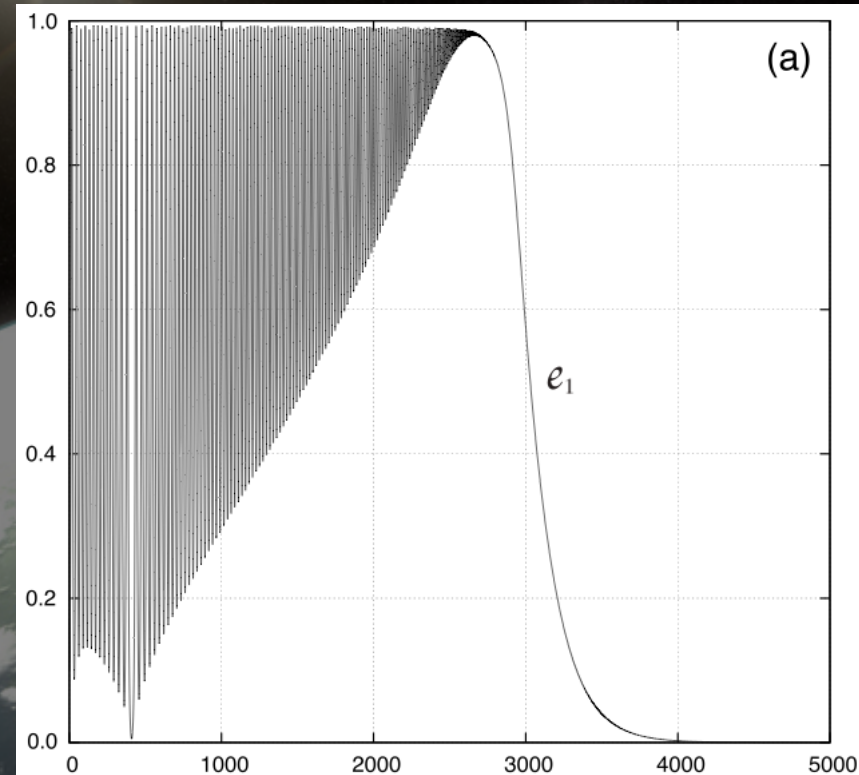
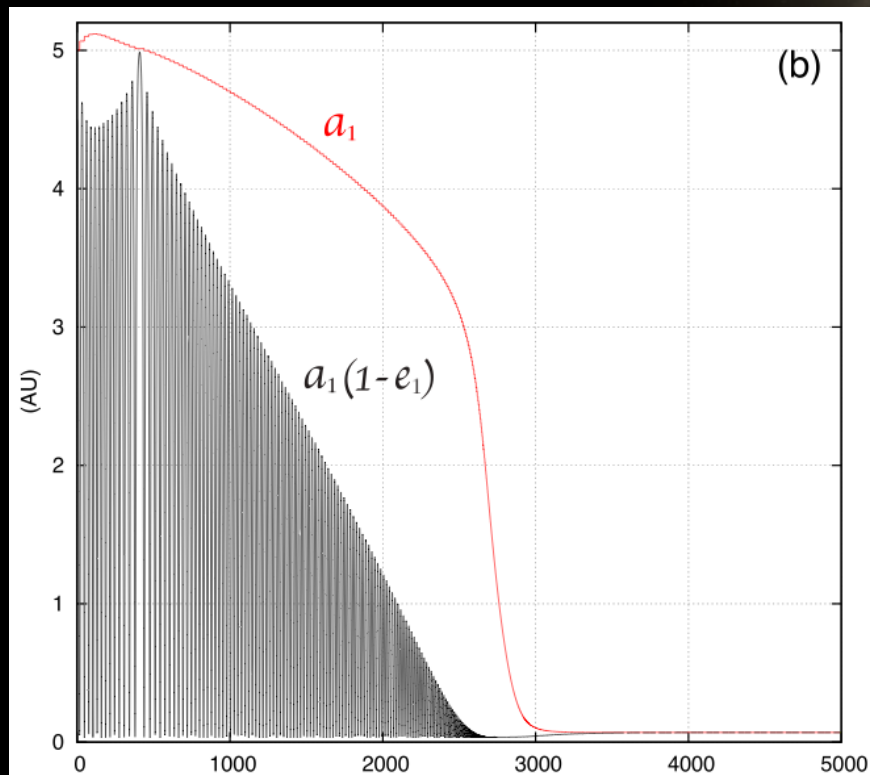


HD 80606

$(a_p = 0.45 \text{ AU}, e_p = 0.92)$

semi-major axis

eccentricity



Wu & Murray (2002)
Frabrycky & Tremaine (2007)
Correia *et al.* (2011)

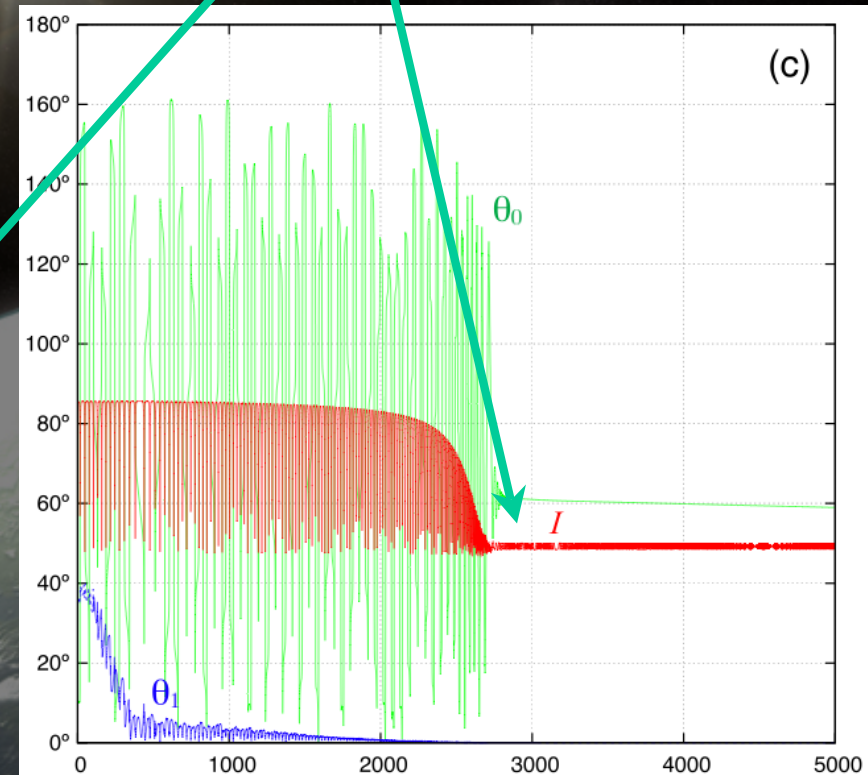
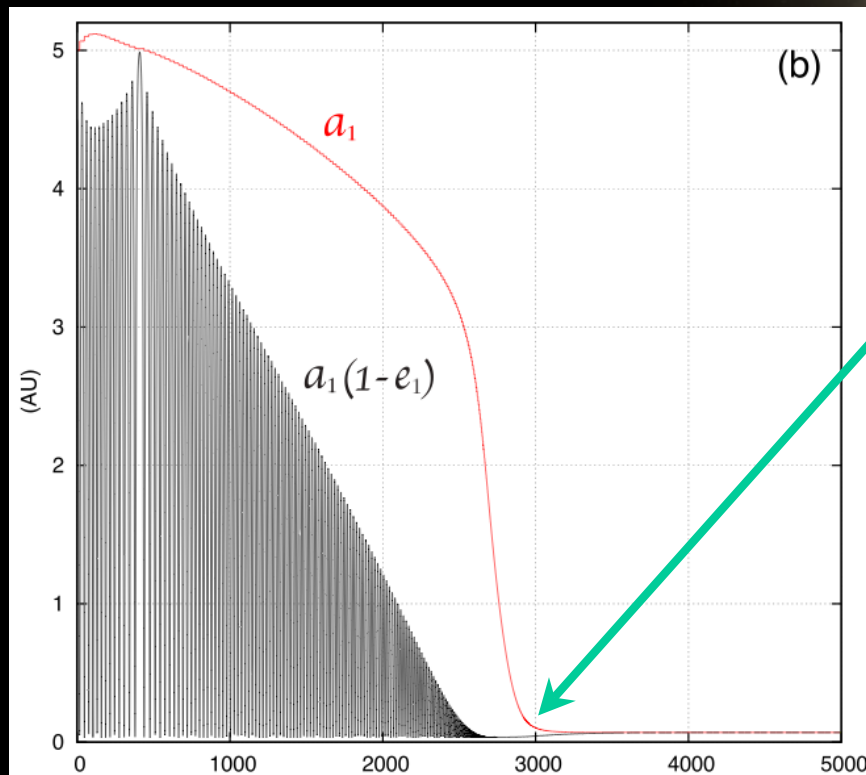
time (Myr)

HD 80606

$(a_p = 0.45 \text{ AU}, e_p = 0.92)$

semi-major axis

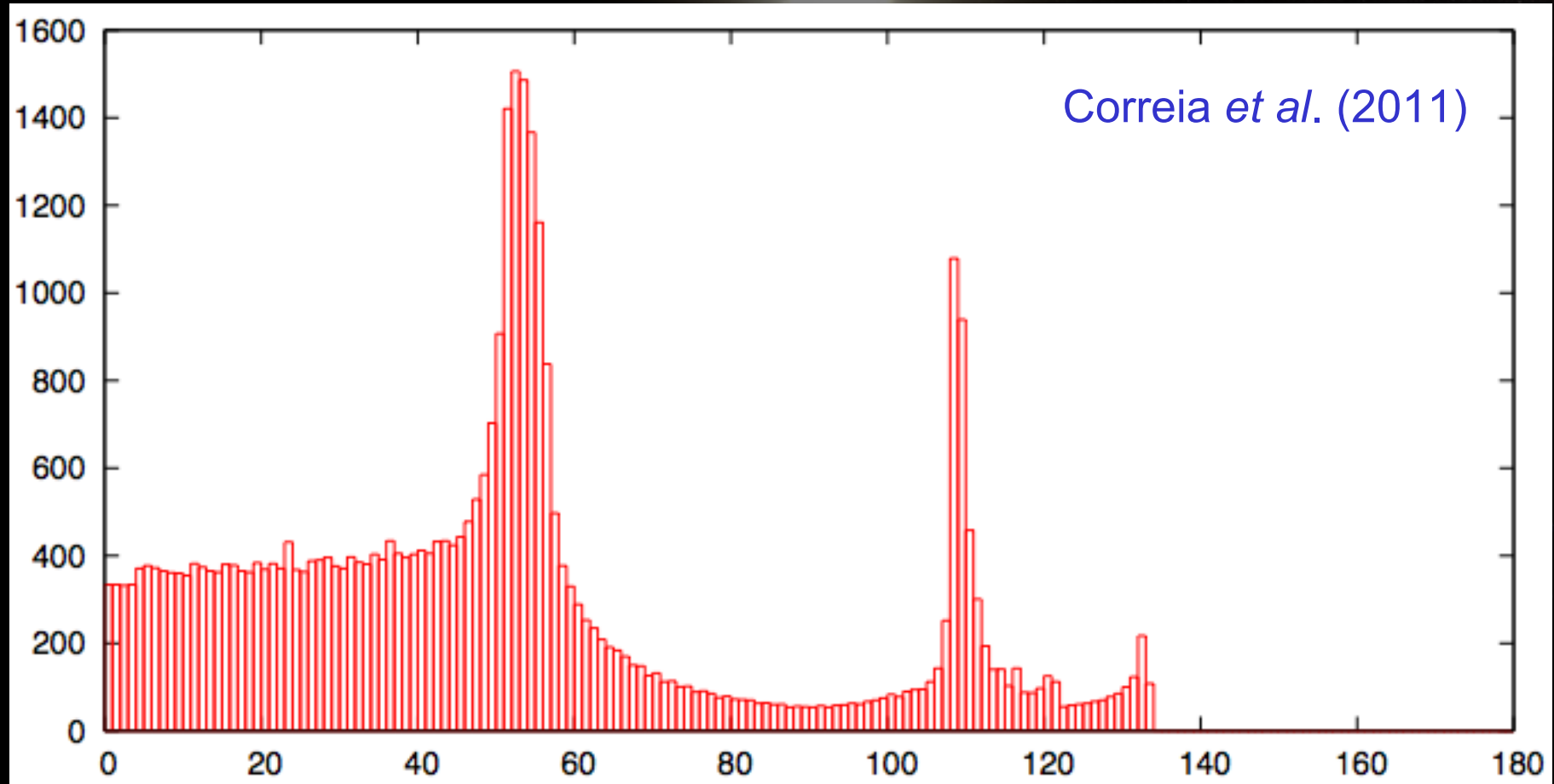
inclination / obliquity



Wu & Murray (2002)
Frabrycky & Tremaine (2007)
Correia *et al.* (2011)

time (Myr)

HD 80606 - misalignment



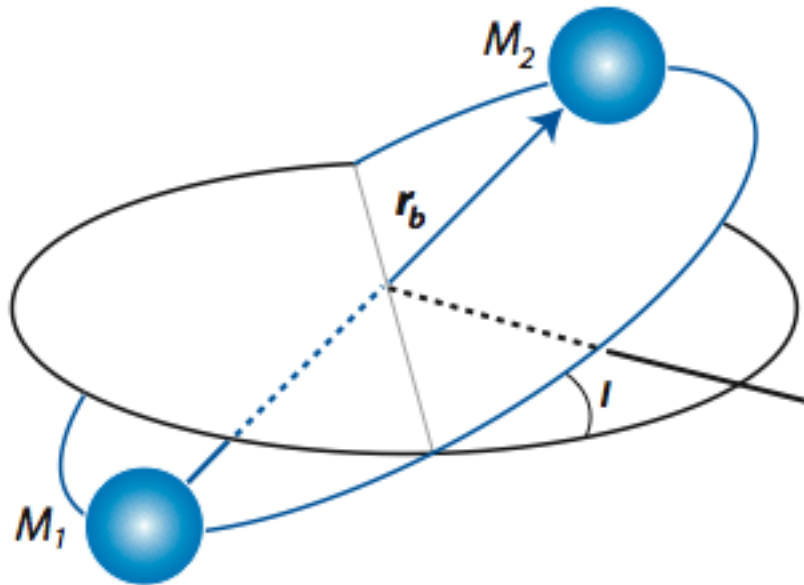
Rossiter-McLaughlin median projected angle = 50°
(Pont *et al.* 2009)

HD 98800

Torres et al. 1995

Boden et al. 2005

$$a_b = 0.98 \text{ AU}, e_b = 0.79, \\ M_1 = 0.7 M_{\text{sun}}, M_2 = 0.6 M_{\text{Sun}}$$



It has a large infrared excess attributed to a circumbinary disc. Substantial extinction suggests that the pair is observed through some of this material.

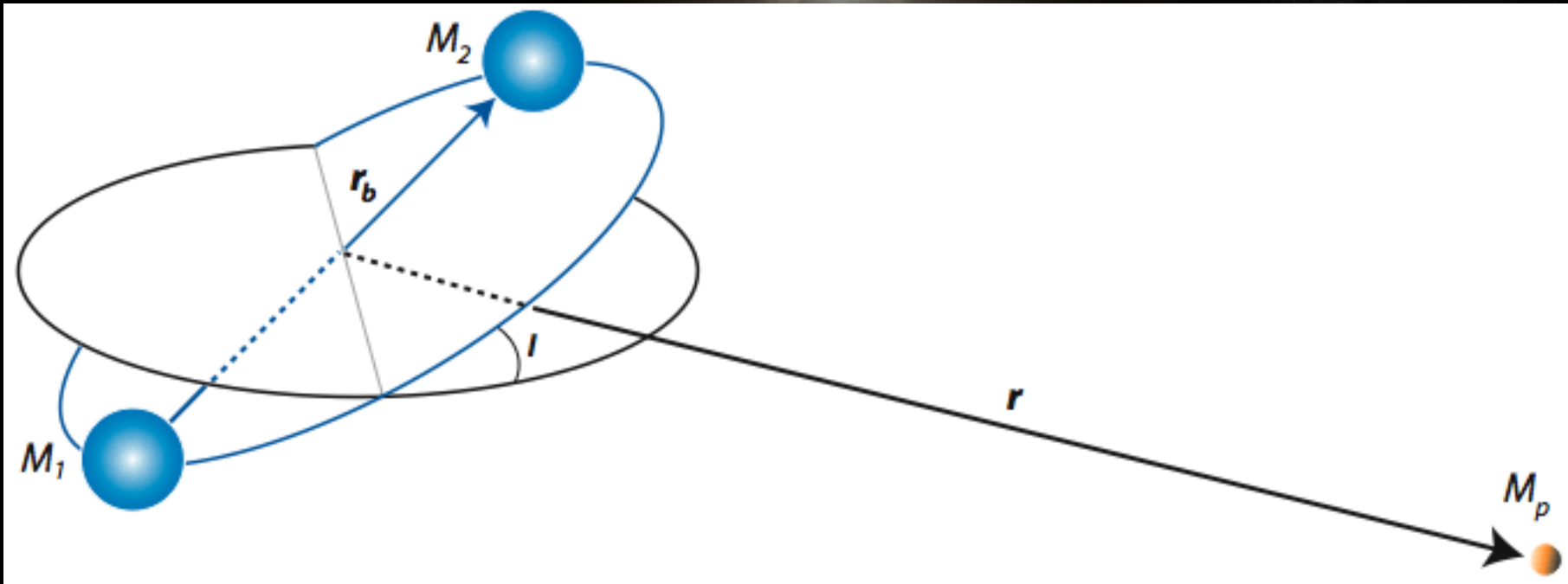
Modeling indicate that this is a T Tauri transition disc that is just reaching the debris disc stage, with a collisional cascade having been recently initiated...

HD 98800

Torres et al. 1995

Boden et al. 2005

$$a_b = 0.98 \text{ AU}, e_b = 0.79, \\ M_1 = 0.7 M_{\text{sun}}, M_2 = 0.6 M_{\text{Sun}}$$



$$i \sim 23^\circ$$

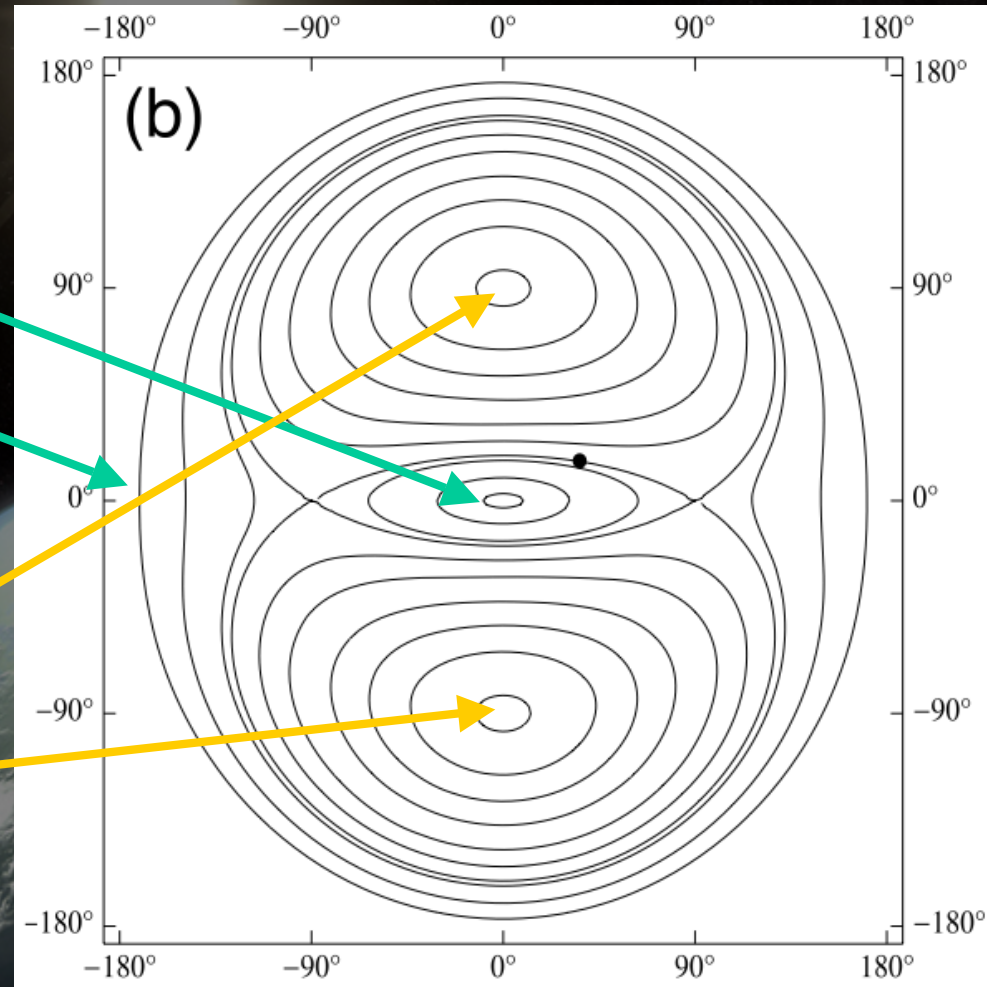
$$a_p = 5.2 \text{ AU}, e_p = 0.5, M_p = 1 M_{\text{Jup}}$$

HD 98800

$$e = 0.79$$

L_2 precesses
around L_1
(or $-L_1$)

L_2 precesses
around e_1
(or $-e_1$)



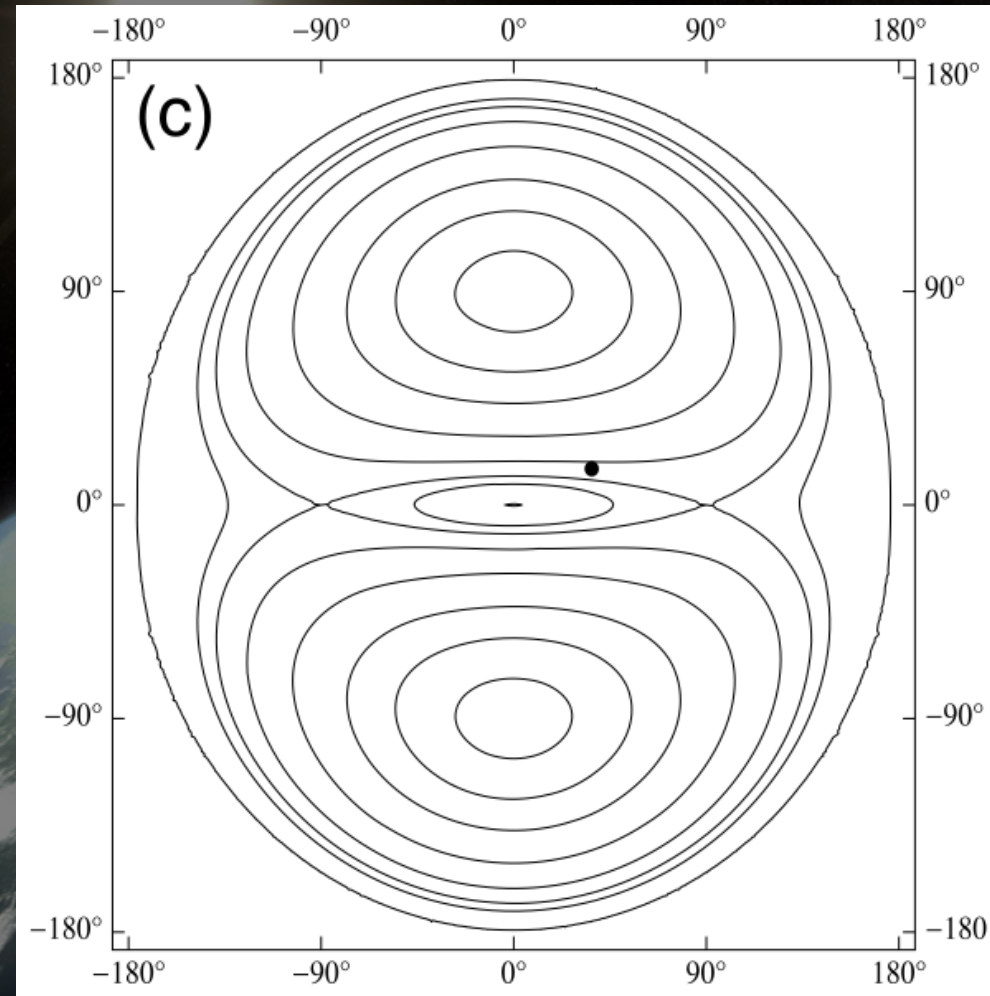
(Farago & Laskar, 2010)

$$(l \sin \omega, l \cos \omega)$$

HD 98800

$$e = 0.9$$

L_2 precesses
around e_1

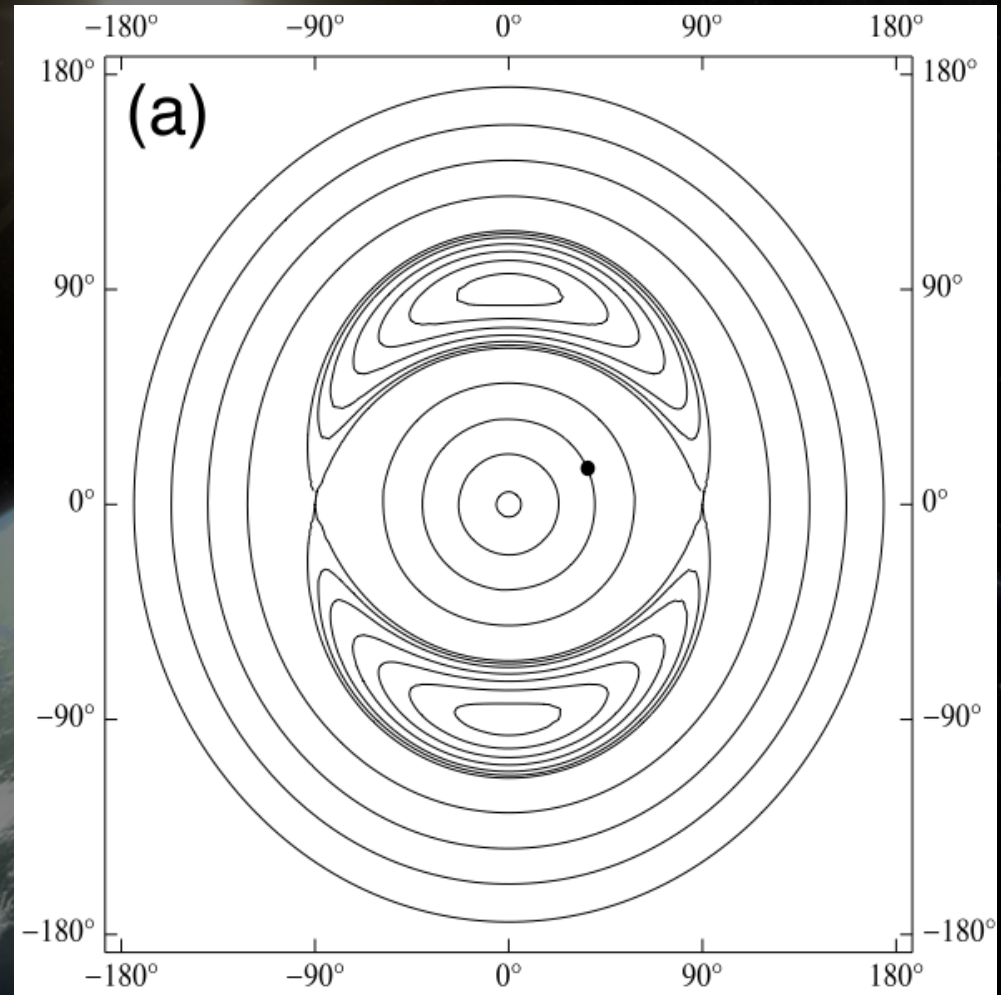


$$(l \sin \omega, l \cos \omega)$$

HD 98800

$$e = 0.5$$

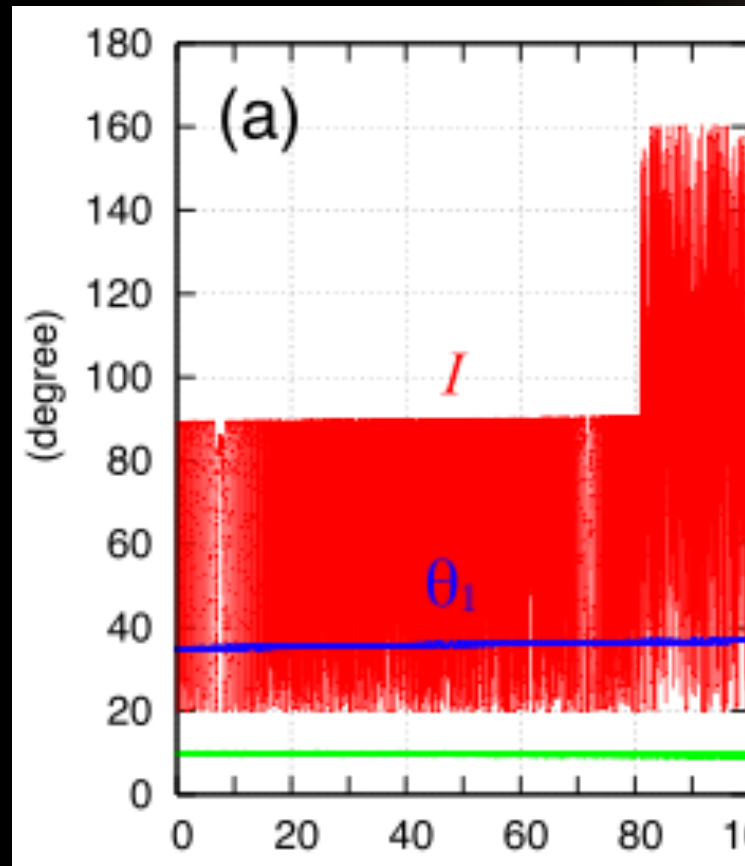
L_2 precesses
around L_1



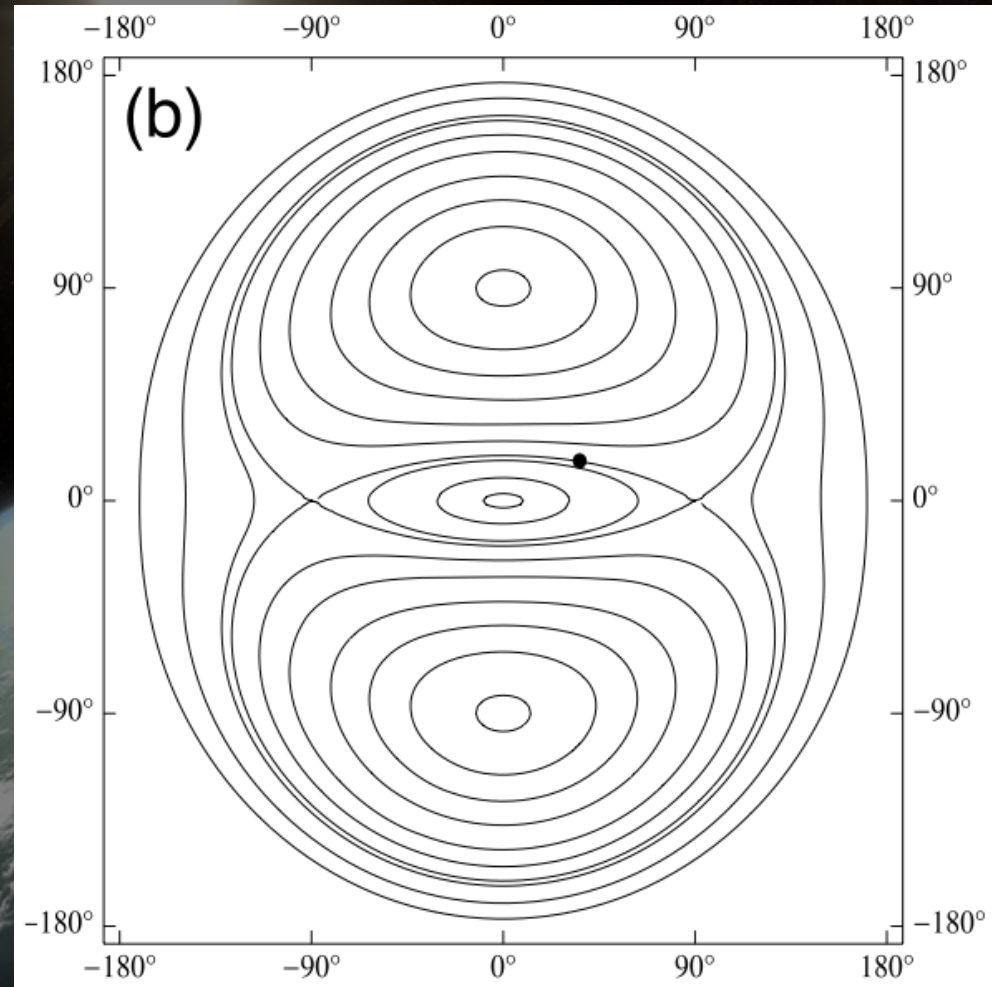
$$(l \sin \omega, l \cos \omega)$$

HD 98800

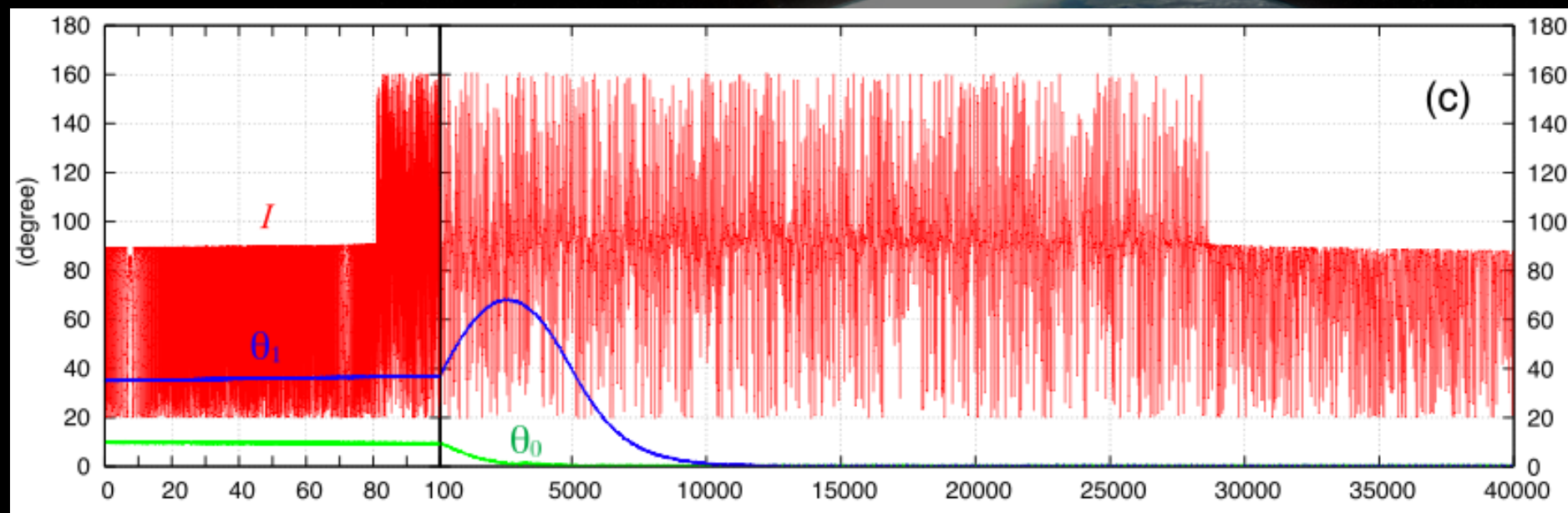
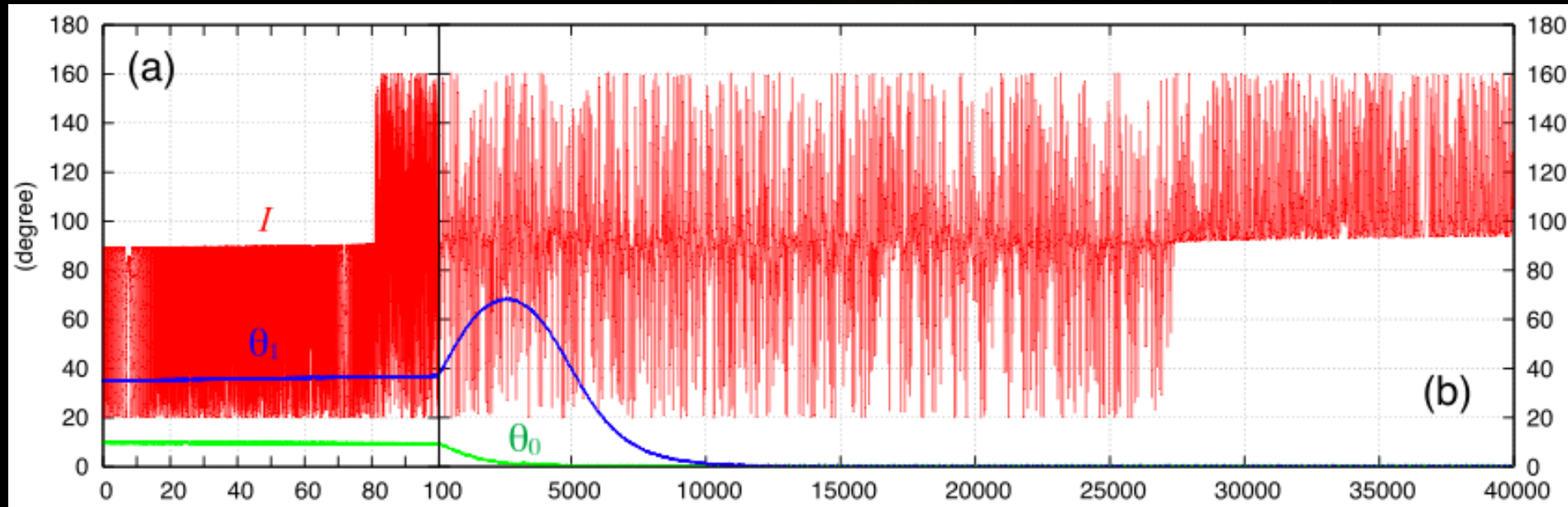
$$e = 0.79$$



time (Myr)



($l \sin \omega, l \cos \omega$)



time (Myr)

Conclusions

- Tidal effects are responsible for a slow secular evolution of the spins and orbits of close-in exoplanets.
- Tidal effects alone align the spin axis, synchronize the rotation and orbital periods, and damp the eccentricity of the orbit.
- Tidal effects combined with planetary perturbations may present unexpected behaviors and lead to stationary equilibria for the spin and the orbits.
- In the restricted inner problem, a body may migrate by means of a tidal-Kozai mechanism. In the restricted outer problem, initial prograde orbits may become retrograde and vice-versa.